

Seminars at FEFU, Vladivostok

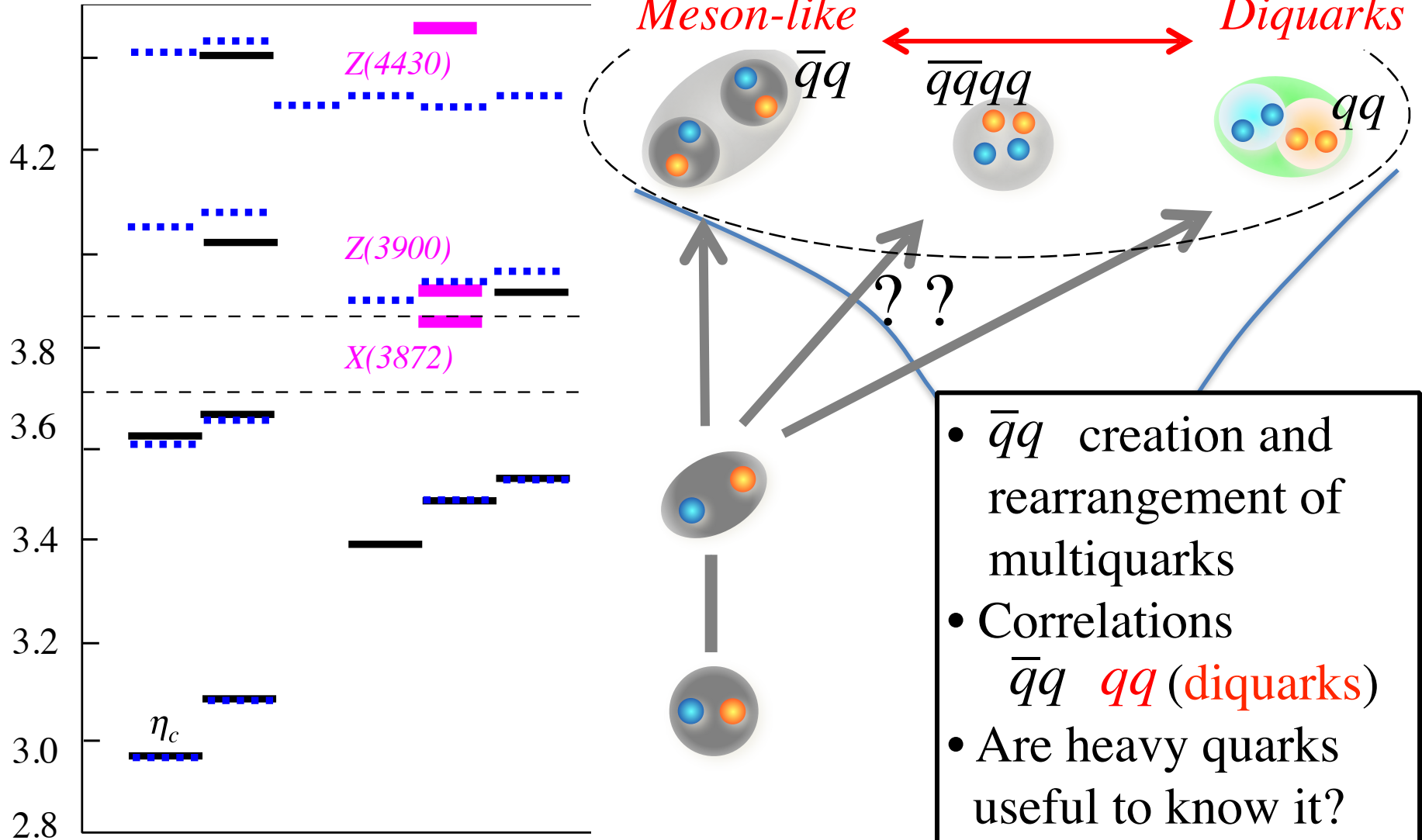
II. Charmed baryons

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Osaka Univ

Quarkonium like states and excitations

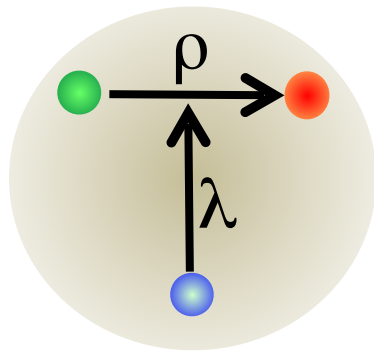


1. ρ , λ modes and diquark motions

A heavy quark may disentangle the fundamental modes λ and ρ

→ place to look at diquark correlations

Isotope-shift: Copley-Isgur-Karl, PRD20, 768 (1979)

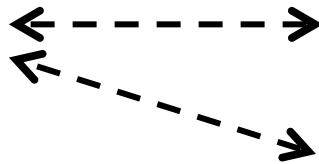


$m_Q = m_{u,d}$

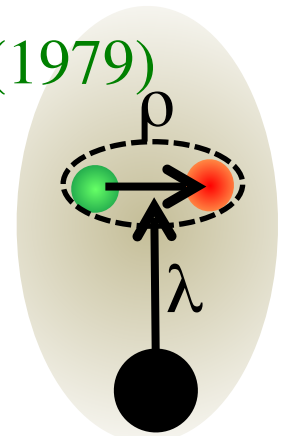
Degenerate

April 20, 2015

$\lambda = \rho$



$\lambda < \rho$



$m_Q \rightarrow \infty$

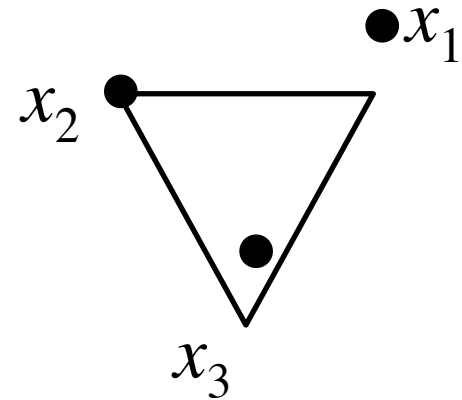
Distinguished

Mixing of λ and ρ

λ excitation $<$ ρ excitation

Harmonic oscillator

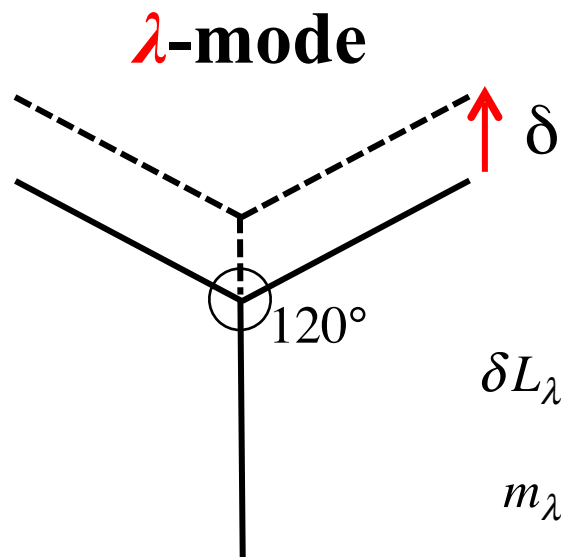
$$H = \frac{p_1^2}{2m} + \frac{p_2^2}{2m} + \frac{p_3^2}{2M} + \frac{k}{2} \left((x_1 - x_2)^2 + (x_2 - x_3)^2 + (x_3 - x_1)^2 \right)$$
$$= \frac{p_\rho^2}{2m_\rho} + \frac{p_\lambda^2}{2m_\lambda} + \frac{k_\rho \rho^2}{2} + \frac{k_\lambda \lambda^2}{2}$$



$$m_\rho = \frac{m}{2}, \quad m_\lambda = \frac{2mM}{M+2m}$$
$$k_\rho = \frac{3}{2}k, \quad k_\lambda = 2k$$

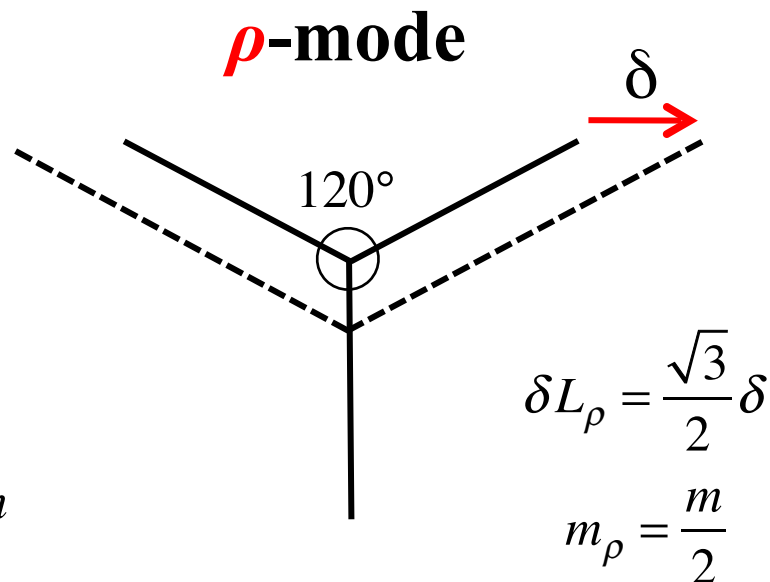
$$\omega_\lambda = \sqrt{\frac{M+2m}{M}}\omega < \omega_\rho = \sqrt{3}\omega$$

String: $V \sim \sigma L$



$$\delta L_\lambda = \delta$$

$$m_\lambda = \frac{2mM}{M+2m} \rightarrow 2m$$



$$\delta L_\rho = \frac{\sqrt{3}}{2} \delta$$

$$m_\rho = \frac{m}{2}$$

$$\omega_\lambda \rightarrow \left(\frac{1}{2}\right)^{1/3} \left(\frac{\sigma^2}{m}\right)^{1/3} < \omega_\rho = \left(\frac{\sigma^2}{m}\right)^{1/3} \rightarrow \left(\frac{3}{2}\right)^{1/3} \left(\frac{\sigma^2}{m}\right)^{1/3}$$

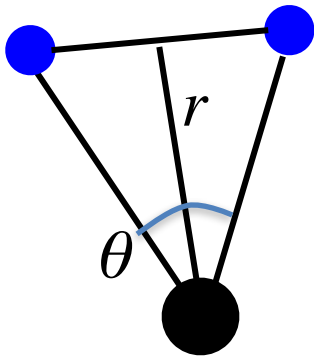
Diquark correlations

Alexandrou, deForcrand, Lucini
PRL 97, 222002 (2006)

$$C_{\Gamma}(\mathbf{r}_u, \mathbf{r}_d, t) = \langle 0 | J_{\Gamma}(\mathbf{0}, 2t) J_0^u(\mathbf{r}_u, t) J_0^d(\mathbf{r}_d, t) J_{\Gamma}^{\dagger}(\mathbf{0}, 0) | 0 \rangle$$

$$J_0^f(\mathbf{r}, t) = :\bar{f}(\mathbf{r}, t) \gamma_0 f(\mathbf{r}, t):, \quad f = u, d,$$

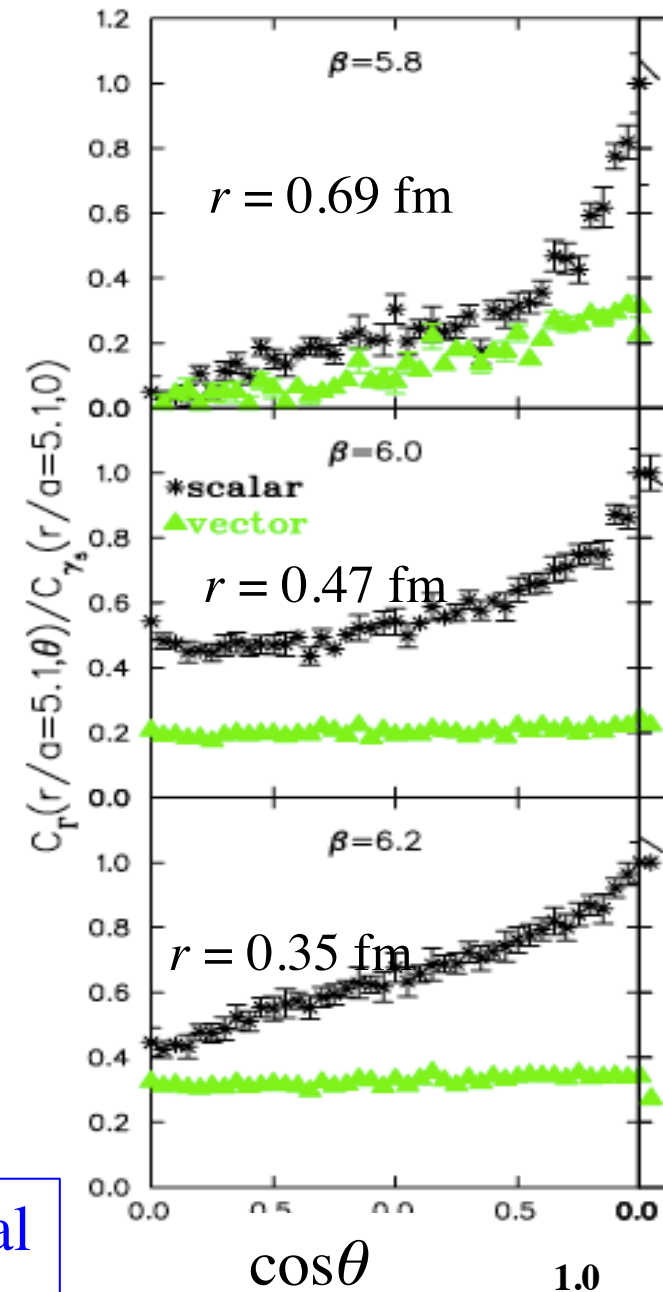
$$J_{\Gamma}(x) = \epsilon^{abc} [u_a^T(x) C \Gamma d_b(x) \pm d_a^T(x) C \Gamma u_b(x)] s_c(x),$$



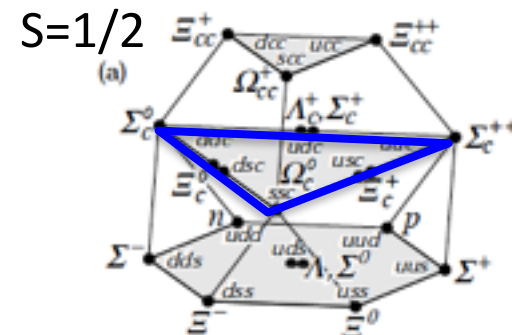
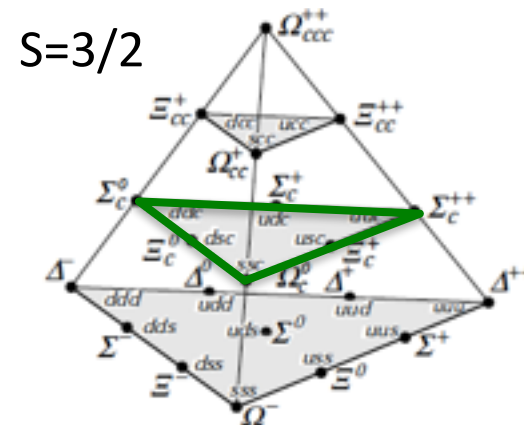
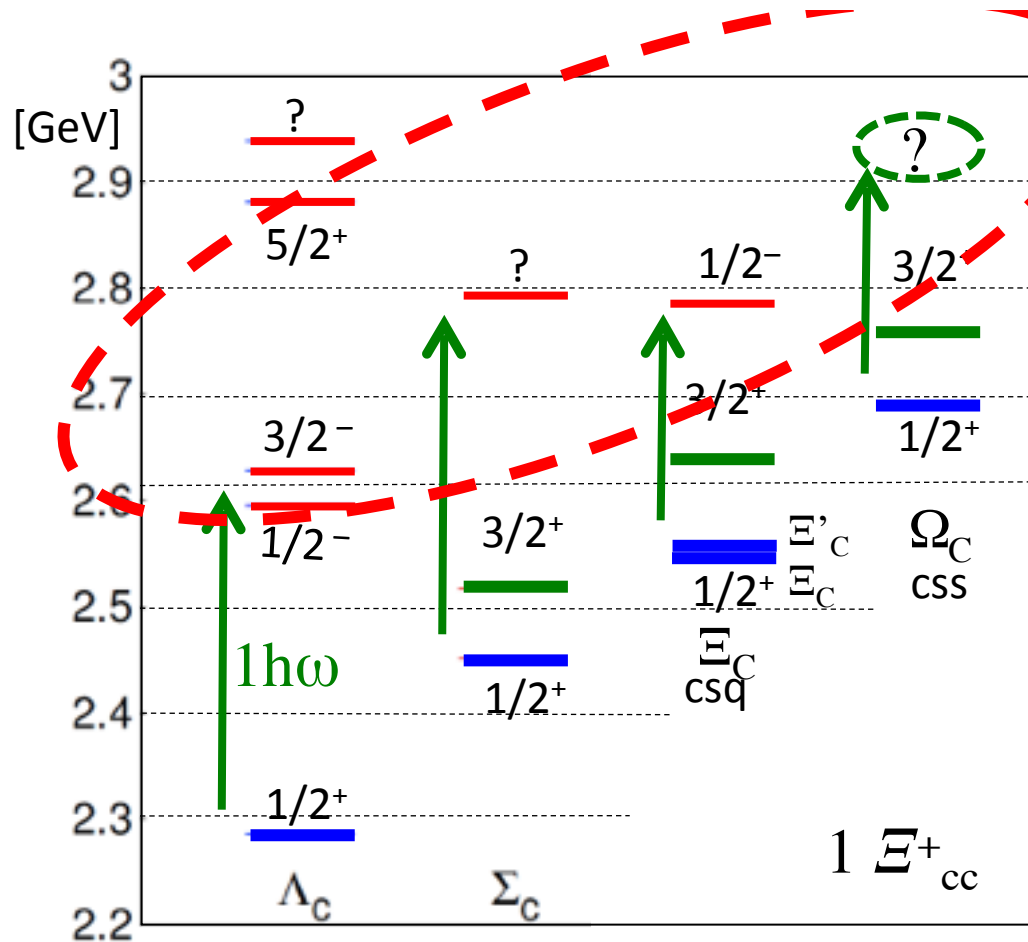
Good diquark
Bad diquark

Indicates significant attraction
between quarks in good diquark pair

See also recent *QCD sum rule study*, Chen et al
arXiv:1502.01103 [hep-ph], To appear in PRD



Spectroscopy: More states at J-PARC



- Excited energies, decay widths are smaller
- Two distinct modes may be different

Quark model calculations

Roberts-Pervin, IJMPA, 23, 2817 (2008)

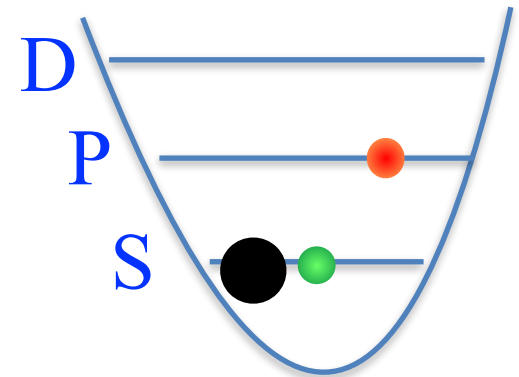
Yoshida, Hiyama, Hosaka, Oka, Phys.Rev. D92 (2015) no.11, 114029

- Hamiltonian

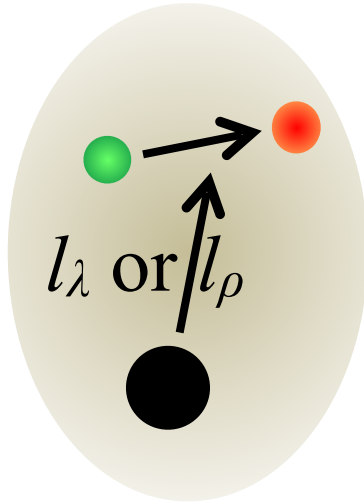
$$H = \frac{p_1^2}{2m_q} + \frac{p_2^2}{2m_q} + \frac{p_3^2}{2M_Q} - \frac{P^2}{2M_{tot}} \\ + V_{conf}(HO) + V_{spin-spin}(Color - magnetic) + \dots$$

- Solved by the **Gaussian expansion** method

- Quark excitations of
P ($l = 1$), D ($l = 2$), ... waves
Many ...



Classifications; p-wave



$$S_{qq} + l = j \text{ (Brown muck)}$$

$$j + 1/2_Q = J_{\text{tot}}$$

$$S_{qq} = 0 \text{ or } 1, \quad l_\lambda \text{ or } l_\rho, \quad j = 0, 1, \dots$$

$$J = 1/2, 3/2, \dots$$

Brown muck

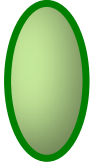
(l_λ, l_ρ)	(l_λ, l_ρ) diquark		J_{tot}	
$(0, 1)$	$(0, 0)$	d^0	1^-	$(1/2, 3/2)^-$
$(0, 0)$	$(0, 1)$	d^1	0^-	$1/2^-$
			1^-	$(1/2, 3/2)^-$
			2^-	$(3/2, 5/2)^-$

← 7 states for Λ and Σ

Only few of them are observed

QM wave functions

λ mode

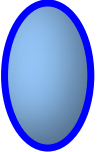


$$\Lambda_c(J^-; \lambda) = \left[[\psi_1(\vec{\lambda})\psi_0(\vec{\rho}), d^0]^1, \chi_c \right]^{J=\frac{1}{2}, \frac{3}{2}} D^0_c$$

charm spin

H.O.(gauss)

flavor



$$\Sigma_c(1/2^+) = \left[[\psi_0(\vec{\lambda})\psi_0(\vec{\rho}), d^1]^1, \chi_c \right]^{\frac{1}{2}} D^1_c$$

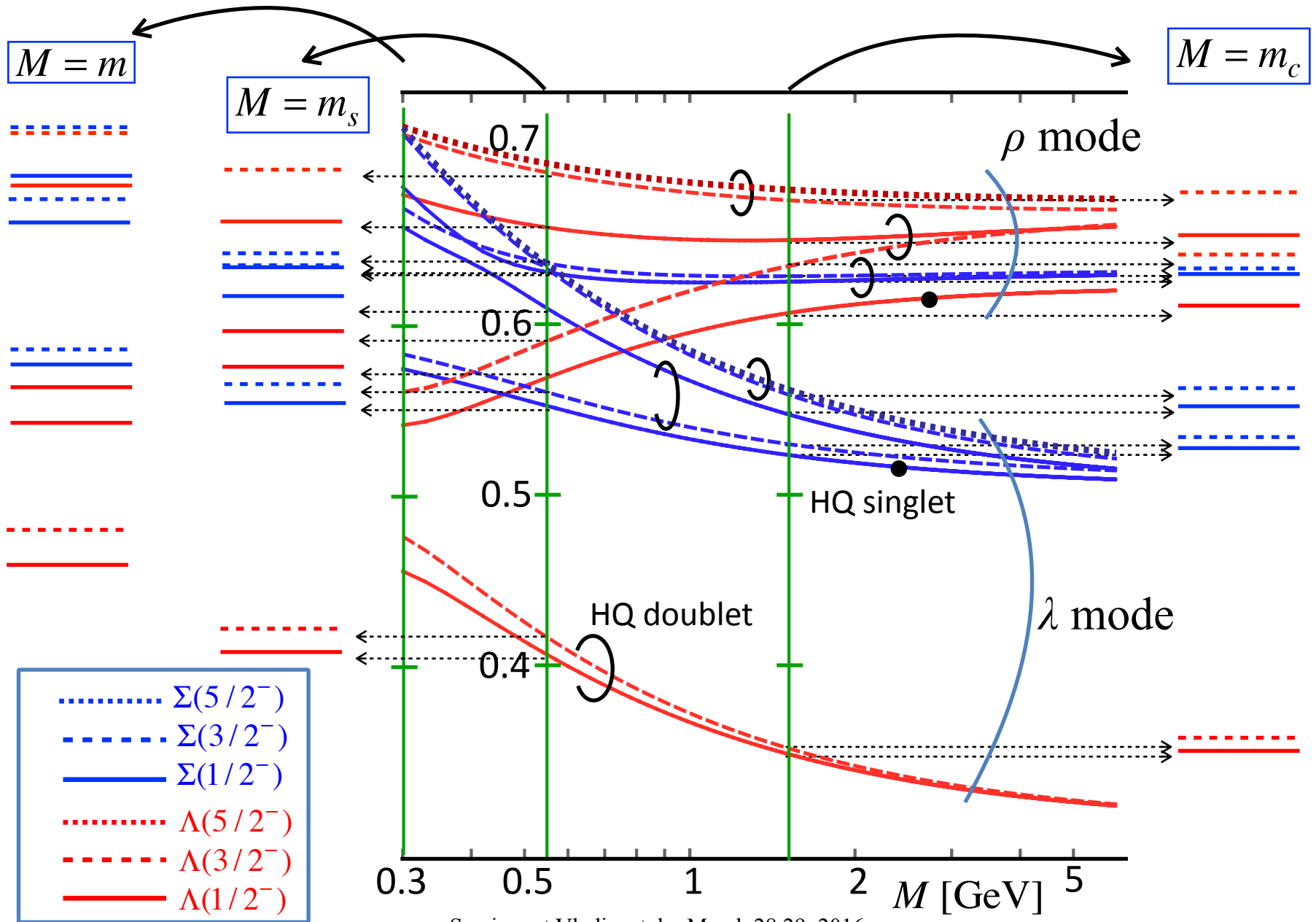
etc.

di-quark spin

Similarly for ρ mode

And mixing of λ and ρ modes

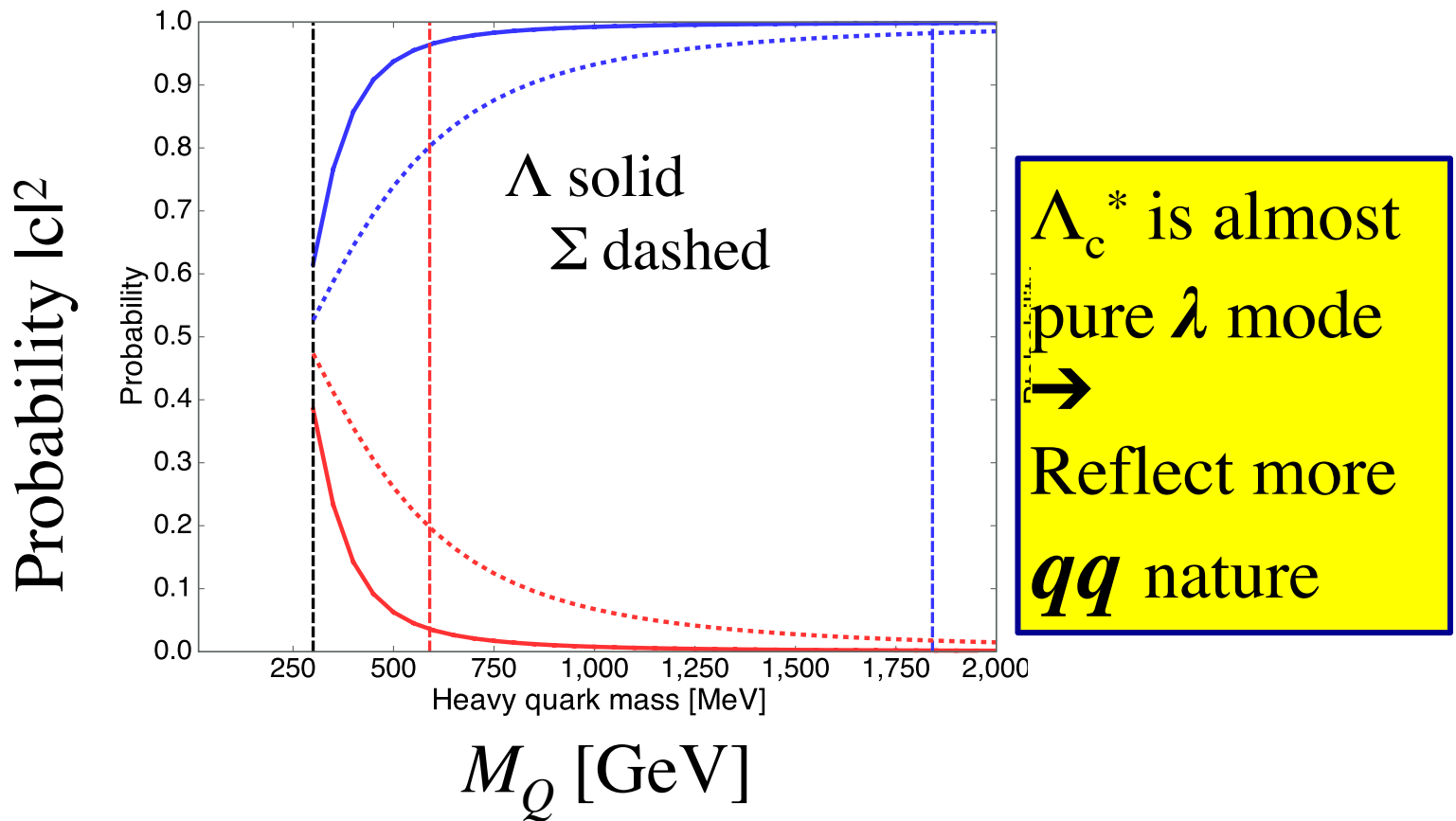
P-wave excitations - $1/2^-$, $3/2^-$



$\lambda\rho$ mixing in the wave functions

Mixing of $\Lambda(\text{phys}) = c_\lambda \Lambda(\lambda) + c_\rho \Lambda(\rho)$

e.g. λ -mode dominant state: How much the other mode mixes?



Decays —Pion emission—

On going, Nagahiro, Yasui, ...

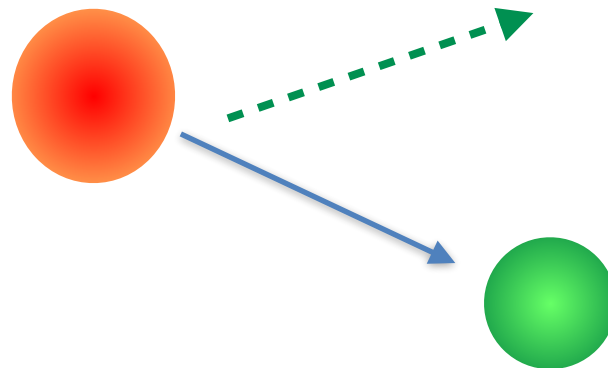
To know the structure,

We study transitions (decay and productions)

Quarks are confined

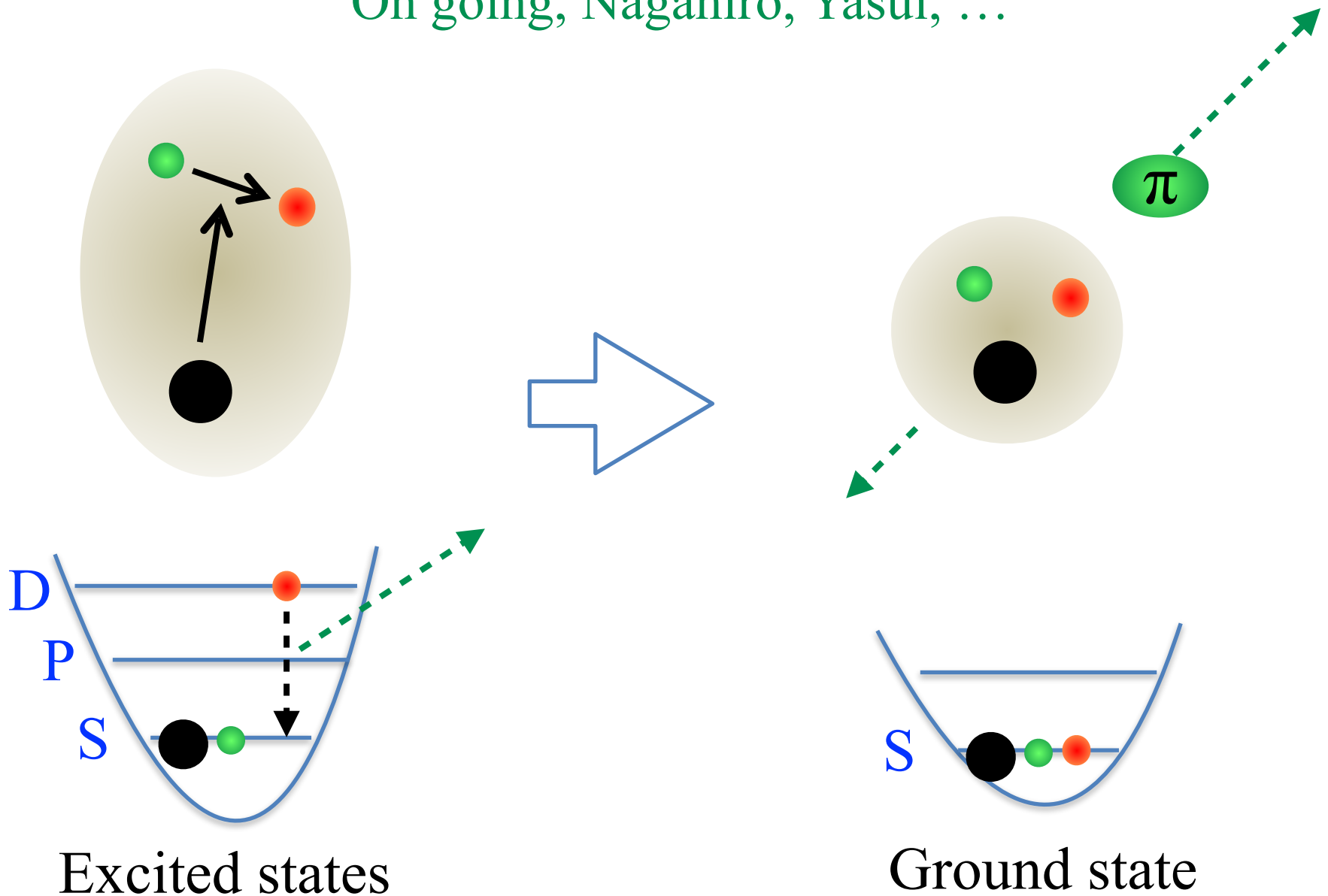
Only transitions through photons, hadrons are available

See the structure that invisible particles form

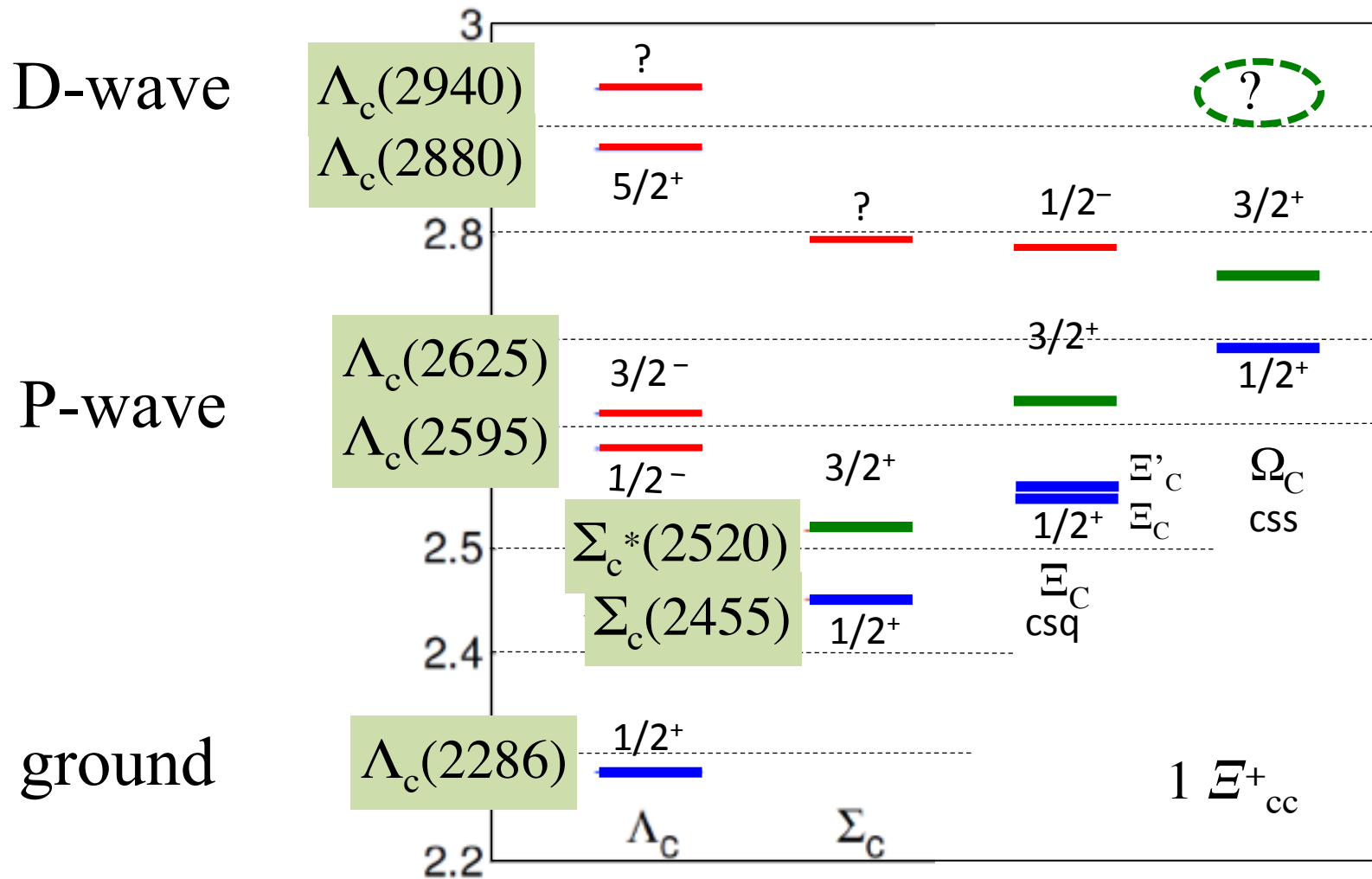


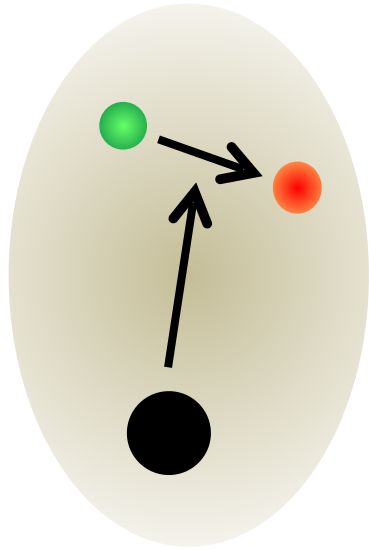
Decays —Pion emission—

On going, Nagahiro, Yasui, ...



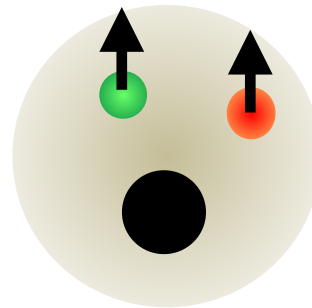
Known spectrum so far



Λ^* 

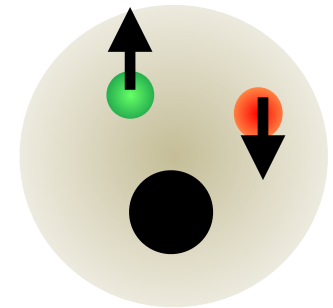
good diquark
 $[[ud]_{S=0, I=0} c]$

$$l = 1$$

 Σ, Σ^* 

bad diquark
 $[[uu]_{S=1, I=1} c]$

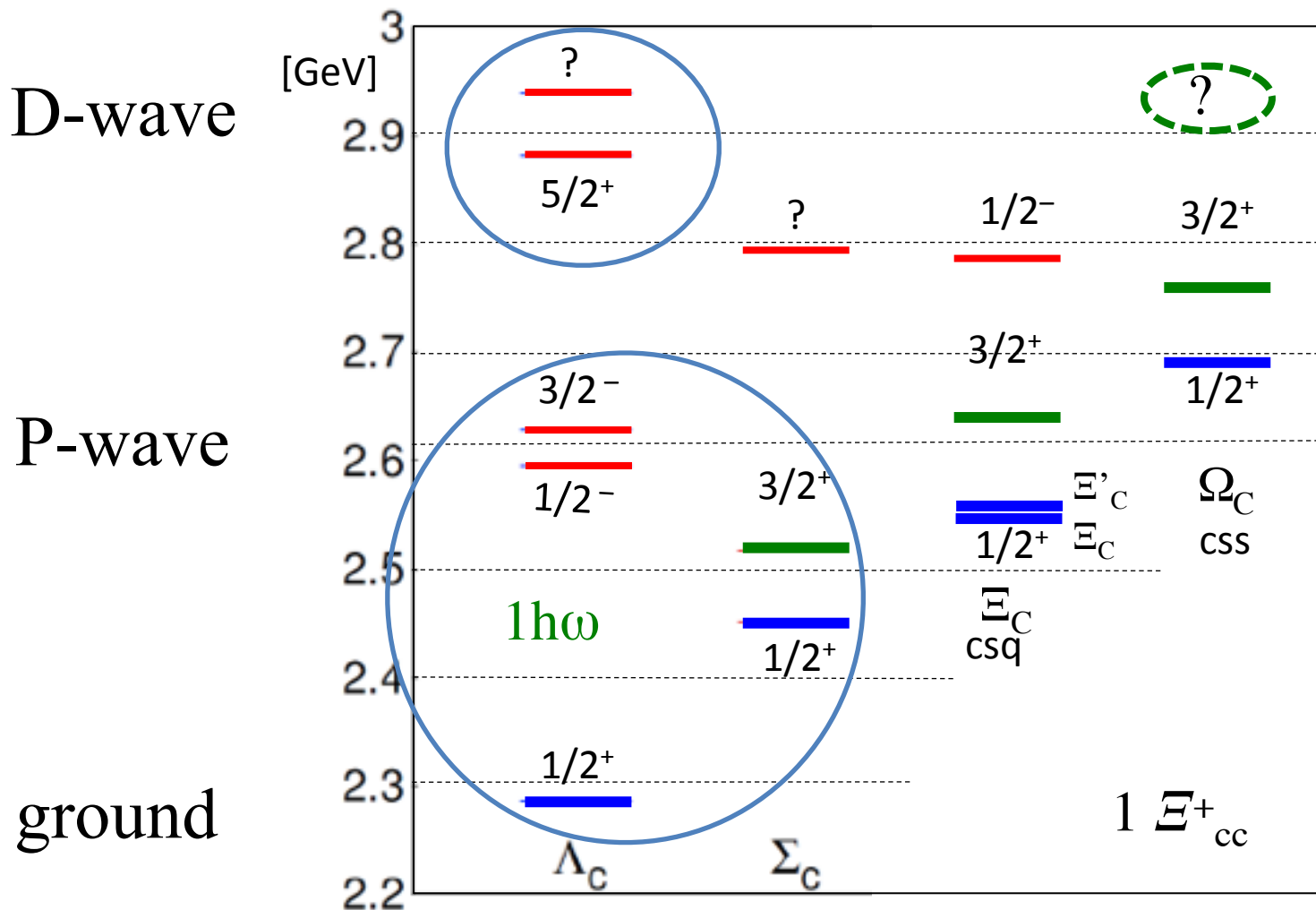
$$l = 0$$

 Λ 

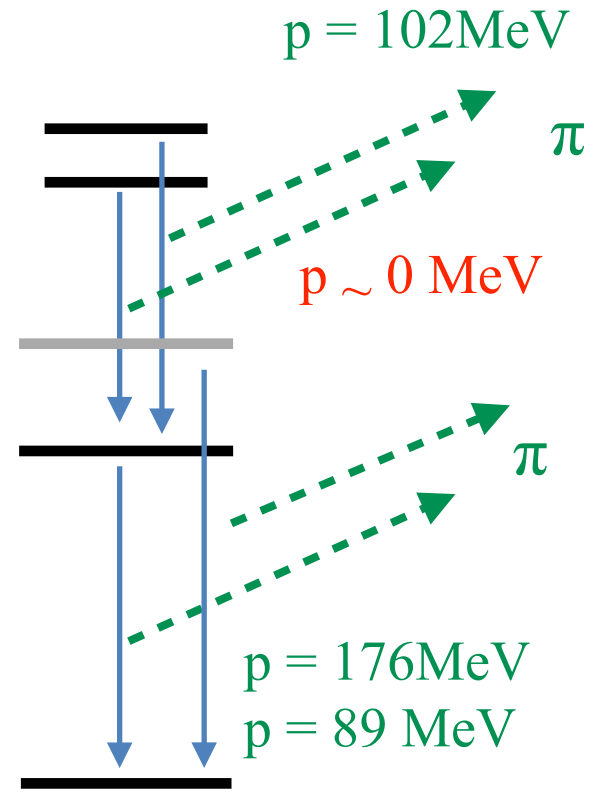
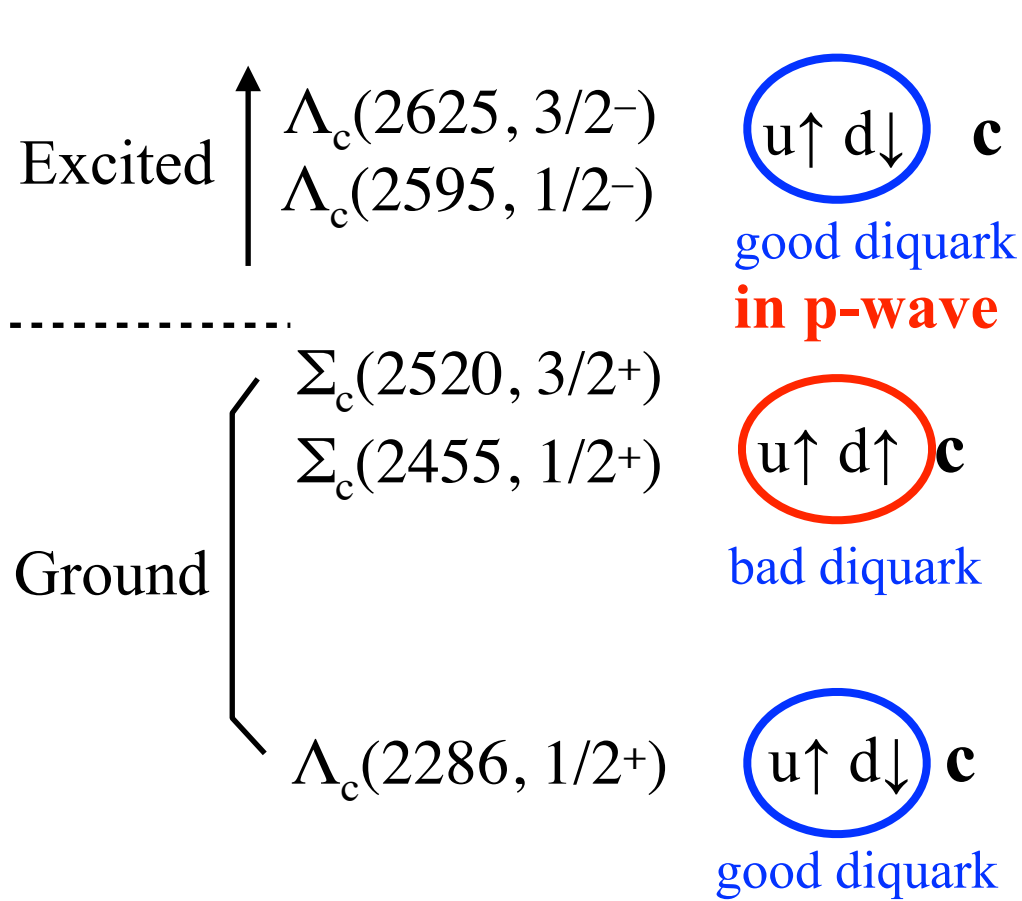
good diquark
 $[[ud]_{S=0, I=0} c]$

$$l = 0$$

Known spectrum so far



Transitions between lower states



low energy
cf: $p(\Delta^- \rightarrow \pi N) = 229 \text{ MeV}$

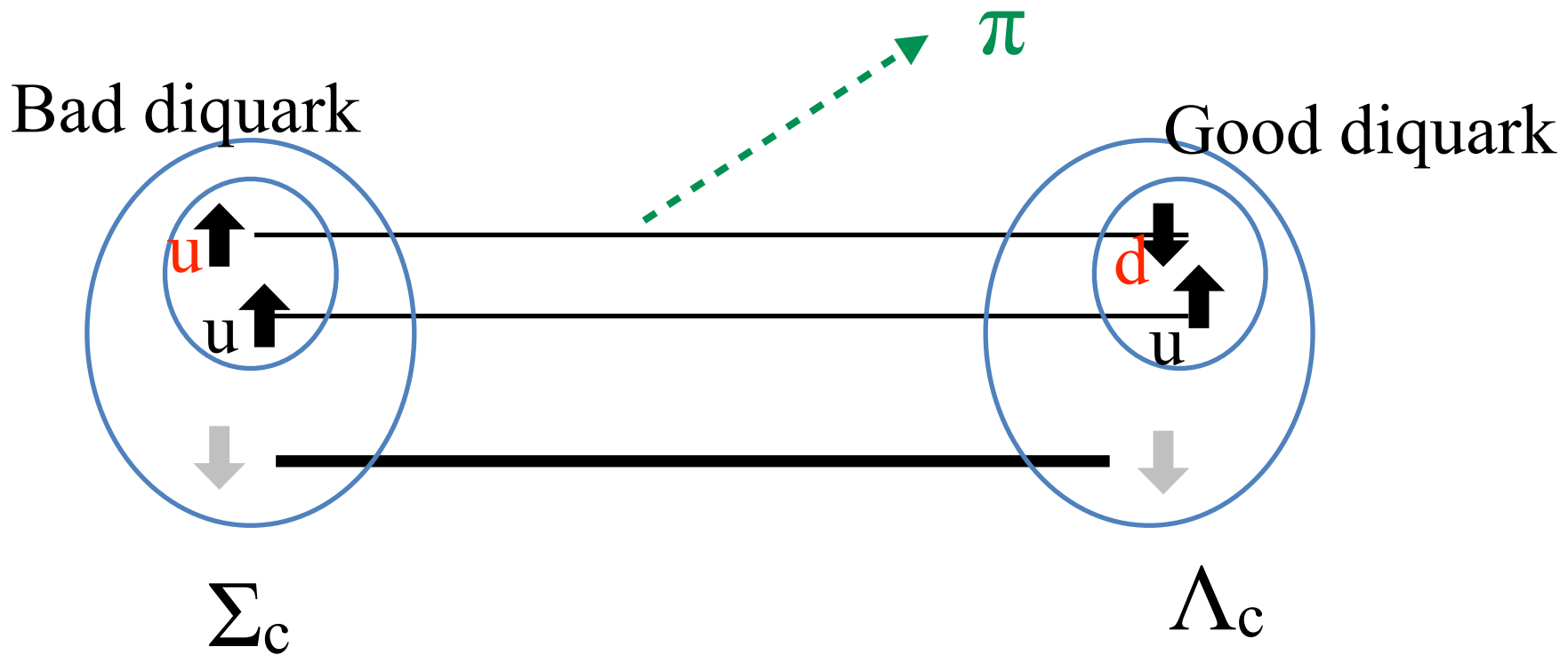
Pion coupling

- Place to look at the *two independent* operators

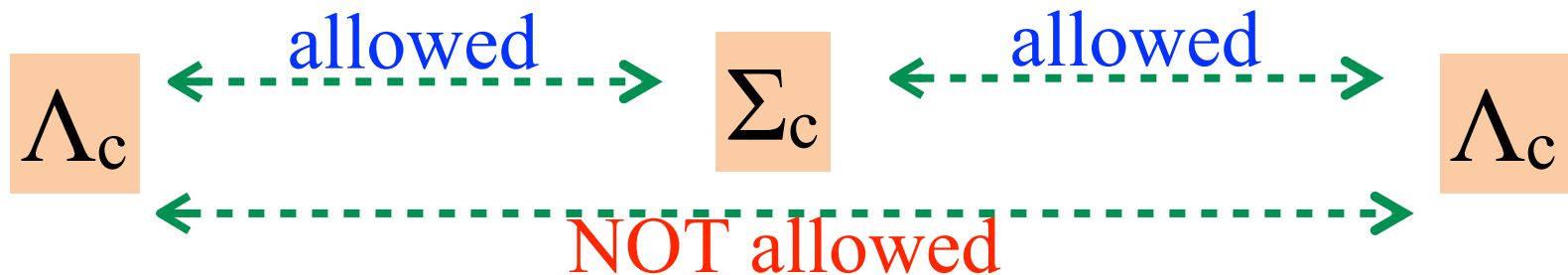
$$\Lambda_c^* \left\{ \begin{array}{c} \vec{p}_i \text{ --- } \vec{p}_f \\ \text{---} \\ \text{---} \end{array} \right\} \Sigma_c \quad \begin{array}{l} \vec{\sigma} \cdot \vec{q} \\ \vec{\sigma} \cdot \vec{p}_i, \vec{\sigma} \cdot \vec{p}_f \\ \bar{q} \gamma_5 q \phi_\pi, \bar{q} \gamma^\mu \gamma_5 q \partial_\mu \phi_\pi \end{array}$$

Chiral dynamics of the NG bosons

$$0^-: \sigma \cdot p$$



- spin-isospin flip between diquark
- charm quark is a spectator



Ground \rightarrow ground transitions

$\Lambda_c(2625, 3/2^-)$

$\Lambda_c(2595, 1/2^-)$

$u\uparrow d\downarrow$ **c**

good diquark
in **p-wave**

$\Sigma_c(2520, 3/2^+)$

$\Sigma_c(2455, 1/2^+)$

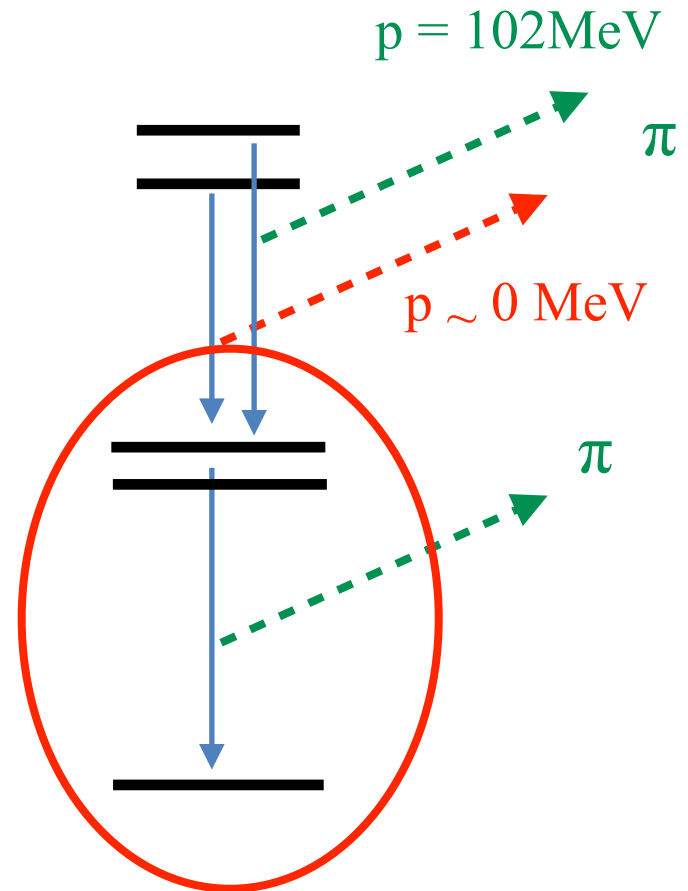
$u\uparrow d\uparrow$ **c**

bad diquark

$\Lambda_c(2286, 1/2^+)$

$u\uparrow d\downarrow$ **c**

good diquark

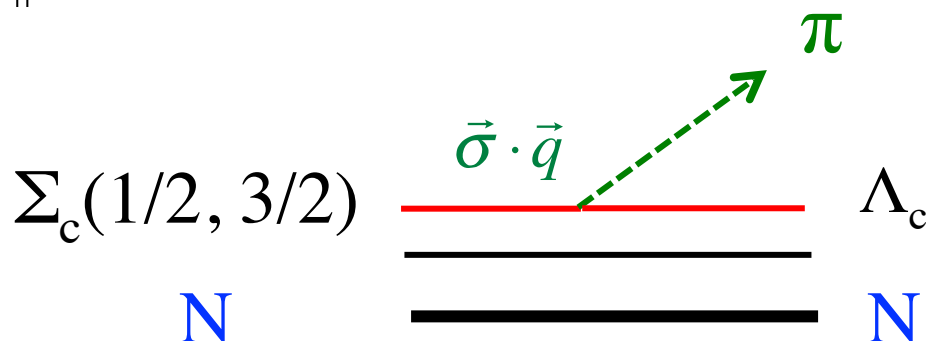


Preliminary results

Ground ($1/2, 3/2^+$) \rightarrow Ground ($1/2^+$)

$B_i J^P$ (MeV)	$\Gamma_{\text{exp}}^{\text{full}}(\Gamma_i)$ (MeV)	q (MeV)	$\Gamma_{\text{th}}(\Sigma_c(J^+)^{++} \rightarrow \Lambda_c^{gs}(1/2^+; 2286)^+ \pi^+)$ (MeV)
$\Sigma_c(2455) 1/2^+$ (2453.98) ($\omega_\pi = 0$ limit)	2.26 (2.26) (2.26)	89	4.27–4.33
$\Sigma_c(2520) 3/2^+$ (2517.9) ($\omega_\pi = 0$ limit)	14.9 (14.9)	176	30.0–31.2

$$g_A^q = 1 \rightarrow g_A^N = 5/3$$



Excited \rightarrow ground transitions

$\Lambda_c(2625, 3/2^-)$

$\Lambda_c(2595, 1/2^-)$

$u\uparrow d\downarrow$ **c**

good diquark
in **p-wave**

$\Sigma_c(2520, 3/2^+)$

$\Sigma_c(2455, 1/2^+)$

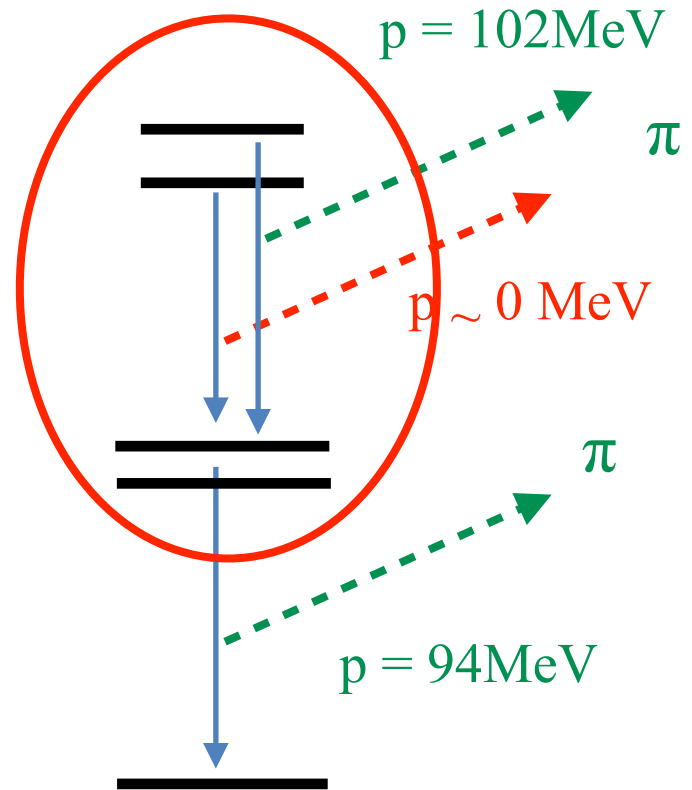
$u\uparrow d\uparrow$ **c**

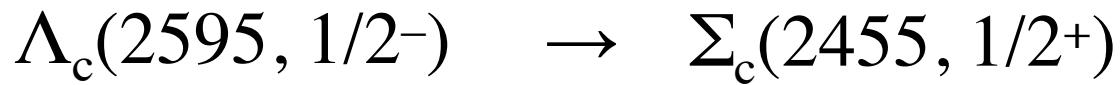
bad diquark

$\Lambda_c(2286, 1/2^+)$

$u\uparrow d\downarrow$ **c**

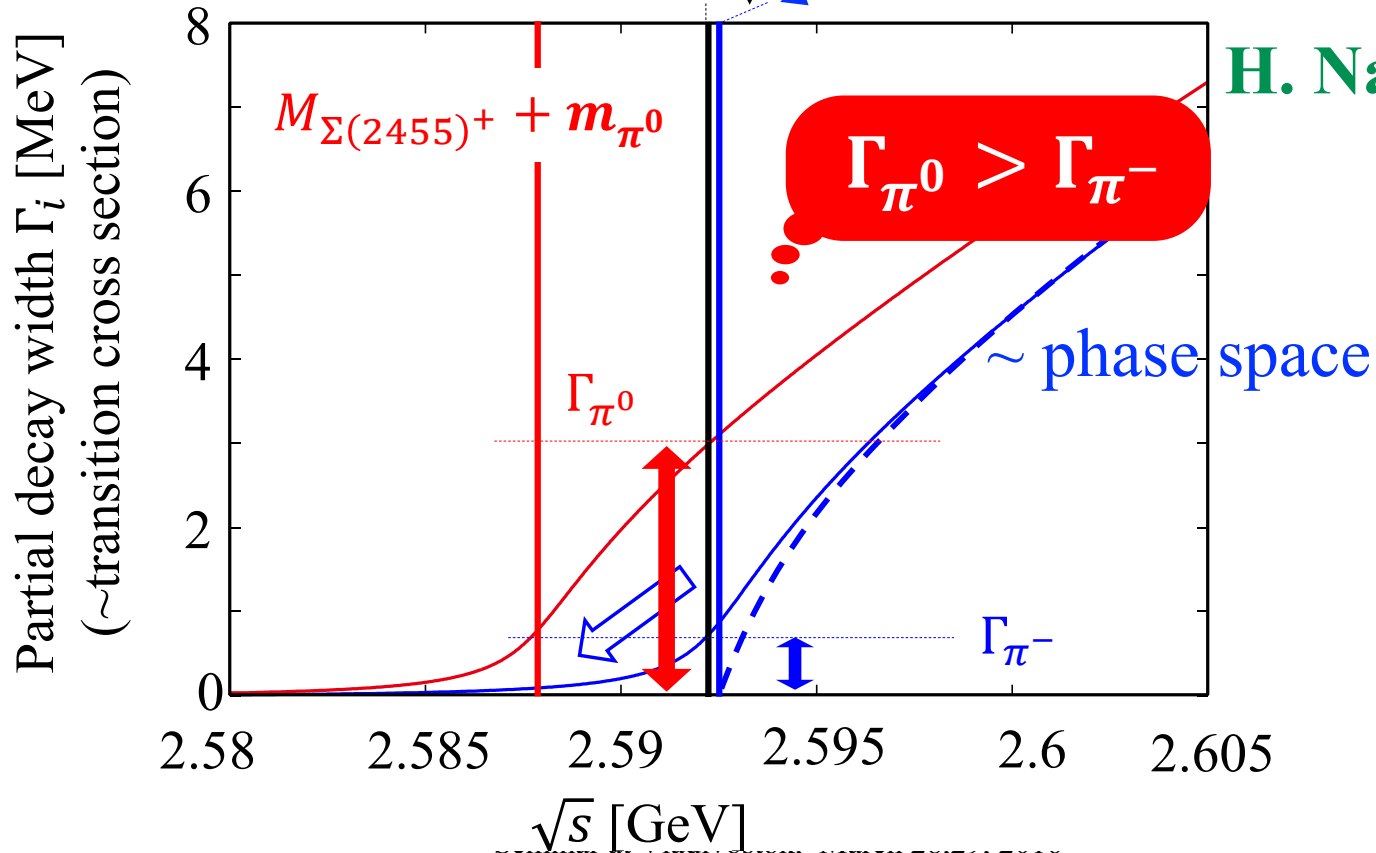
good diquark





Isospin breaking right on the threshold

$$M_{\Lambda(2595)} < M_{\Sigma(2455)^{++}} + m_{\pi^-} \Rightarrow \Gamma_{\pi^-} = 0 ?$$



Preliminary results

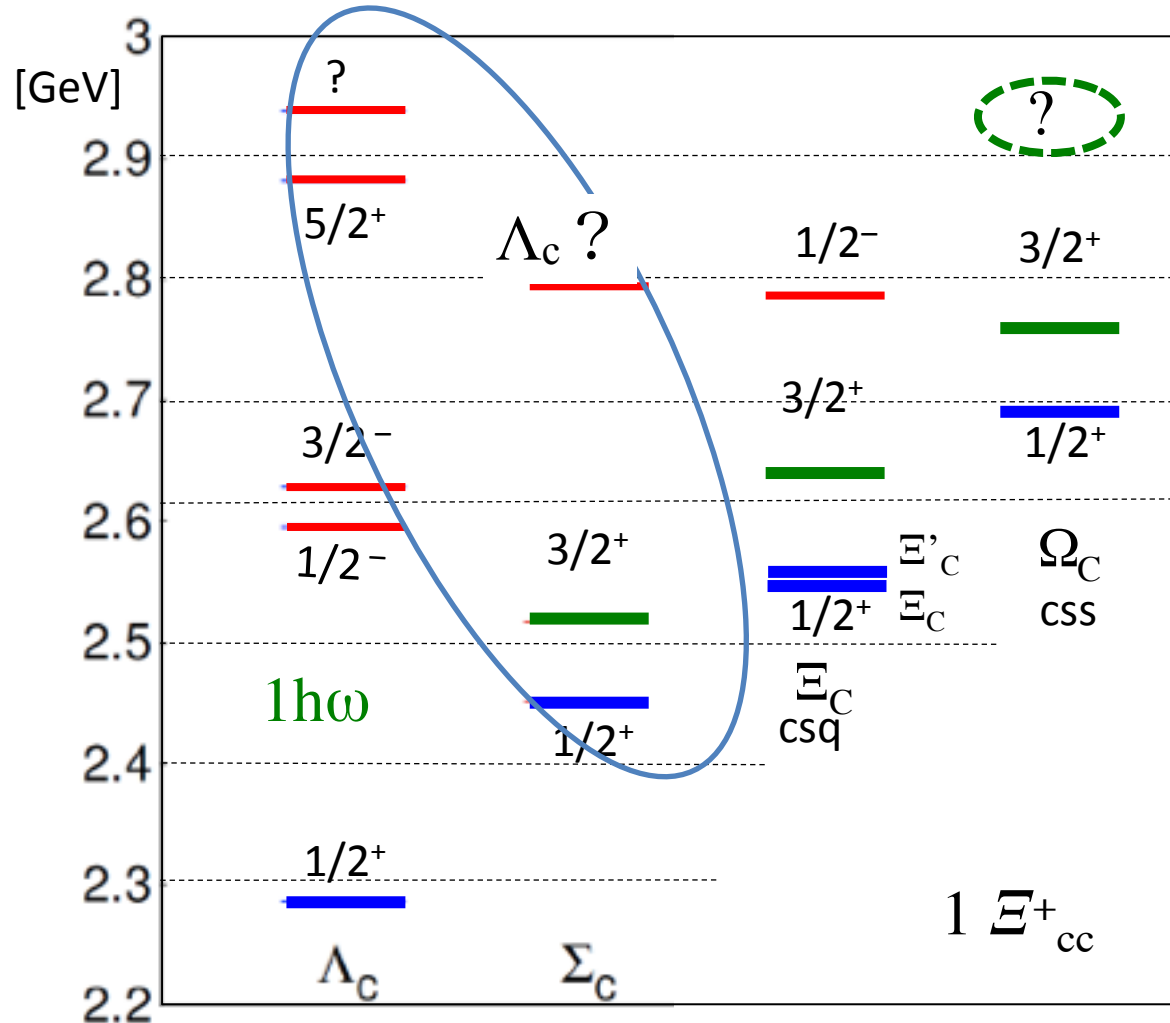
P-wave excitation \rightarrow Ground ($1/2^+$)

			$\Lambda_c(2592; 1/2^-)$	$\Lambda_c(2625; 3/2^-)$
PDG value Γ_{tot} [MeV] (3body含む)			2.6 ± 0.26	< 0.95
Γ^{theo} ($\Sigma_c\pi$) [MeV]	λ -mode	$1/2^-$	1.7 – 3.2	18.3 – 44.4
		$3/2^-$	-	0.075 – 0.14
	ρ -mode	$1/2^- (j = 0)$	0	0
		$1/2^- (j = 1)$	7.3 – 13.1	80 - 188
		$3/2^- (j = 1)$	-	0.036 – 0.66
		$3/2^- (j = 2)$	-	0.078 – 0.12
		$5/2^- (j = 2)$	-	0.030 – 0.053

 **7 states**

Decays —Pion emission—

On going, Nagahiro, Yasui, ...



$\Lambda_c(2765), \Lambda_c(2880), \Lambda_c(2940) \rightarrow \Sigma_c \pi, \Sigma_c^* \pi$ decay widths ?

in PDG, $\text{BR}(\Sigma_c^{(*)} \pi / \text{total})$ are not shown

\rightarrow sum-up $\Sigma_c \pi + \Sigma_c^* \pi$ decay widths and compare to $\Gamma_{\text{tot}}^{\text{PDG}}$

		$\Lambda_c(2765)?^?$	$\Lambda_c(2880)5/2^+$	$\Lambda_c(2940)?^?$
PDG value total Γ [MeV]		50	5.8 ± 1.1	17_{-6}^{+8}
This work $\Gamma_{\Sigma\pi} + \Gamma_{\Sigma^*\pi}$ [MeV]	λ -mode	$1/2^-$	65 – 146	112 – 255
	J^-	$3/2^-$	52 – 104	129 – 249
		$1/2^- (j = 0)$	0	0
		$1/2^- (j = 1)$	325 – 675	503 – 1130
	ρ -mode	$3/2^- (j = 1)$	211 – 414	440 – 921
	J^-	$3/2^- (j = 2)$	9 – 13	53 – 68
		$5/2^- (j = 2)$	6 – 9	42 – 55
		$1/2^+ (j = 0)$	1.6 – 4.4	3.7 – 13.5
	λ -mode	$3/2^+ (j = 2)$	4.6 – 10.8	16.2 – 39.2
	J^+	$5/2^+ (j = 2)$	1.9 – 4.3	11.1 – 26.1

$\Lambda_c(2765), \Lambda_c(2880), \Lambda_c(2940) \rightarrow \Sigma_c \pi, \Sigma_c^* \pi$ decay widths ?

in PDG, $\text{BR}(\Sigma_c^{(*)} \pi / \text{total})$ are not shown

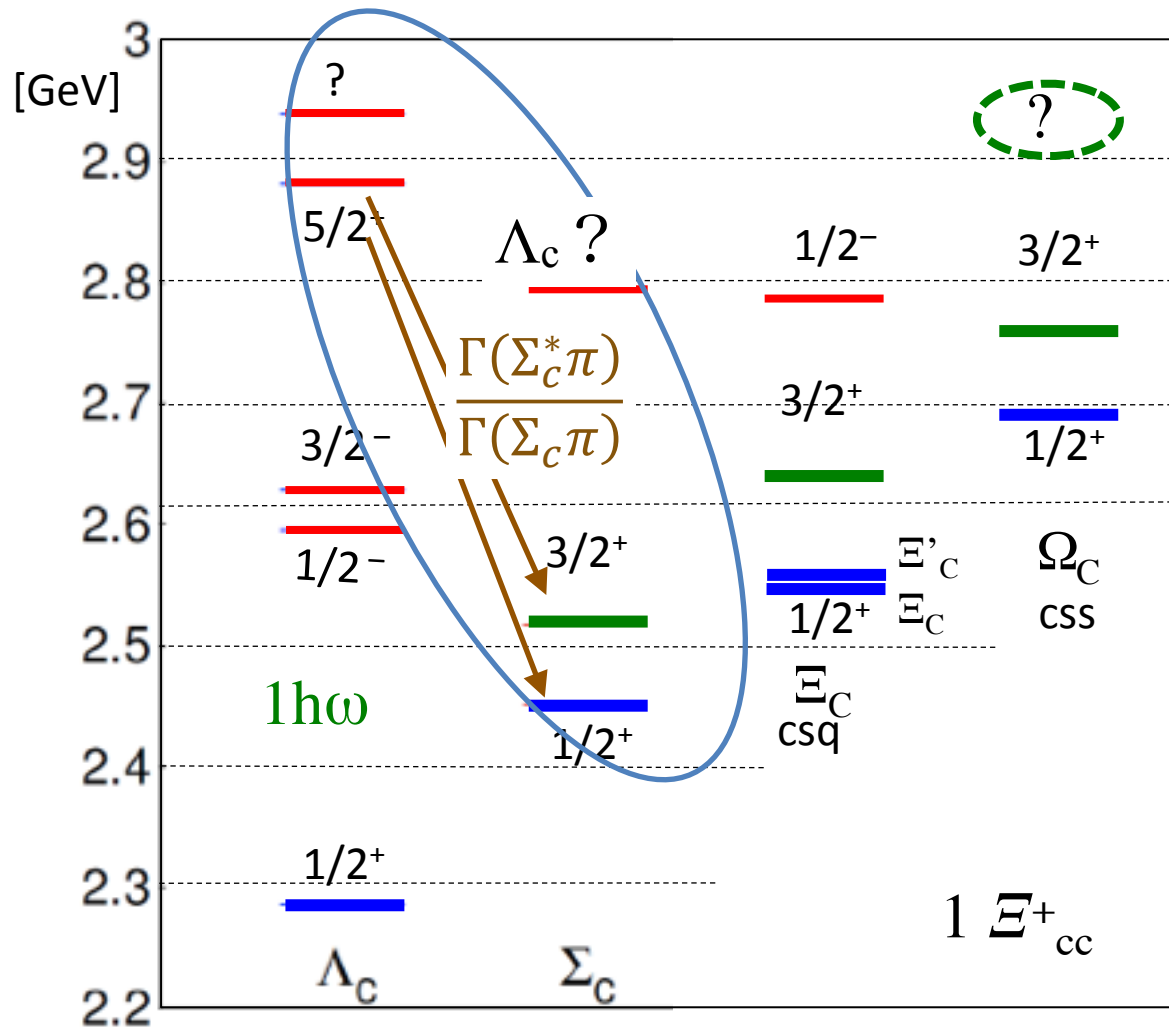
\rightarrow sum-up $\Sigma_c \pi + \Sigma_c^* \pi$ decay widths and compare to $\Gamma_{\text{tot}}^{\text{PDG}}$

		$\Lambda_c(2765)?^?$	$\Lambda_c(2880)5/2^+$	$\Lambda_c(2940)?^?$
PDG value total Γ [MeV]		50	5.8 ± 1.1	17_{-6}^{+8}
This work $\Gamma_{\Sigma_c \pi} + \Gamma_{\Sigma_c^* \pi}$ [MeV]	λ -mode			
	J^-			
	$1/2^-$	65 – 146	112 – 255	145 – 314
	$3/2^-$	52 – 104	129 – 249	182 – 332
	$1/2^- (j = 0)$	0	0	0
	$1/2^- (j = 1)$	325 – 675	503 – 1130	558 – 1301
	$3/2^- (j = 1)$	211 – 414	440 – 921	537 – 1155
	$3/2^- (j = 2)$	9 – 13	53 – 68	96 – 119
	$5/2^- (j = 2)$	6 – 9	42 – 55	80 – 101
	$3/2^+ (j = 0)$	1.6 – 4.4	3.7 – 13.5	3.8 – 17.5
$3/2^+ (j = 2)$	4.6 – 10.8	16.2 – 39.2	24.8 – 61.4	
$5/2^+ (j = 2)$	1.9 – 4.3	11.1 – 26.1	19.8 – 46.5	

$\frac{\Gamma(\Sigma_c^* \pi)}{\Gamma(\Sigma_c \pi)}$ are different

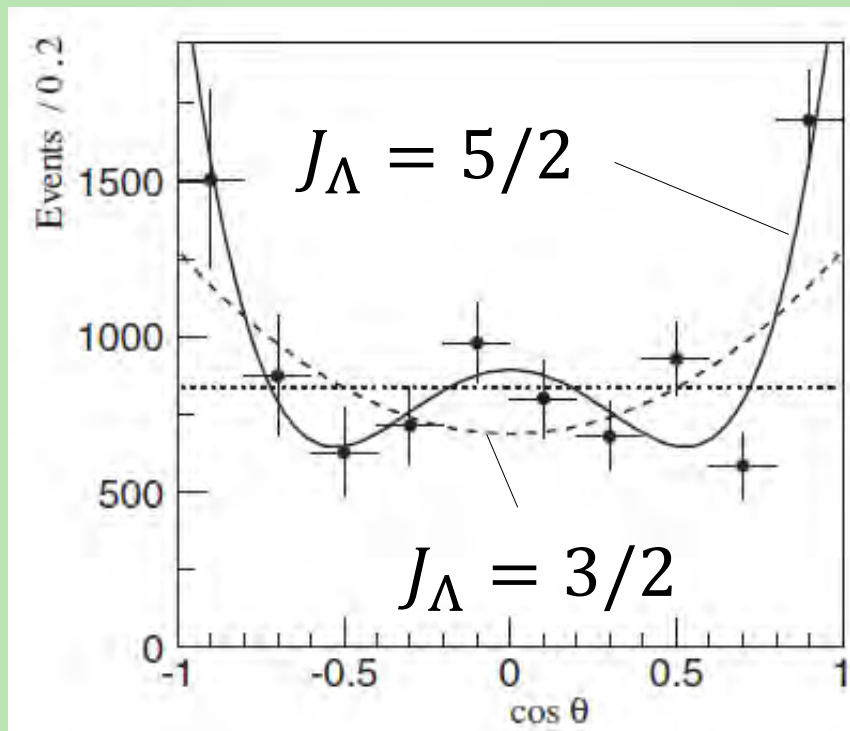
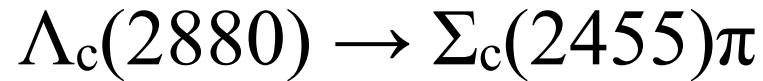
Decays —Pion emission—

On going, Nagahiro, Yasui, ...



Spin and parity of $\Lambda(2880)$

- Spin \leftarrow decay angular distribution



[Mizuk et al., (Belle), PRL98(07)262001, FIG. 3]

- Parity ← decay ratio

$$R = \frac{\Gamma[\Lambda_c(2880, 5/2^P) \rightarrow \Sigma_c^*(2520, 3/2^+) \pi]}{\Gamma[\Lambda_c(2880, 5/2^P) \rightarrow \Sigma_c^*(2455, 1/2^+) \pi]}$$

$$= \frac{\Gamma[F \text{ wave}] + \Gamma[P \text{ wave}]}{\Gamma[F \text{ wave}]}$$

Ignoring P wave and use the HQ symmetry

$$R = \frac{\Gamma[F \text{ wave}] + \cancel{\Gamma[P \text{ wave}]}}{\Gamma[F \text{ wave}]} \sim 0.23_{HQ}$$

$$\sim 0.225 \pm 0.062 \pm 0.025_{EXP}$$

Isgur-Weise, PRL66(91)1130
Cheng et al., PRD75(07)014006

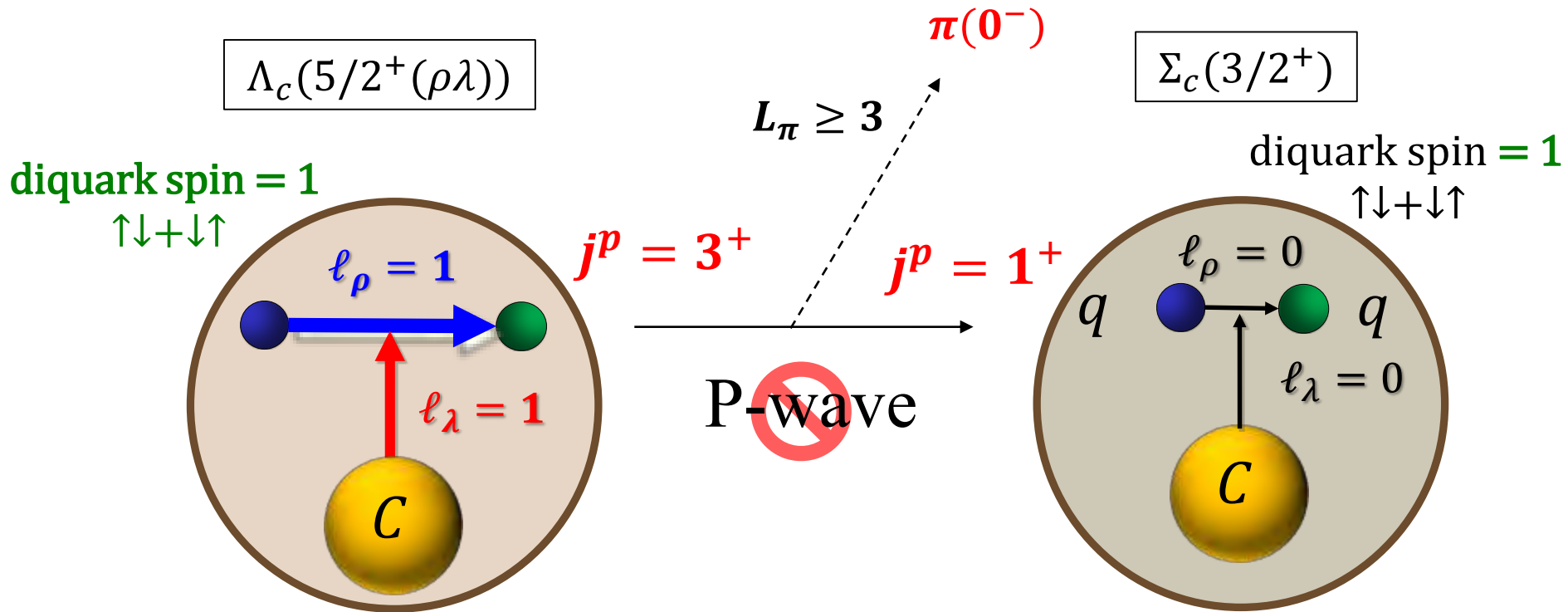
Generally, however, P wave exists and is larger than F wave term → How can we explain data?

j ... Brown muck spin

				$\Lambda_c(2880)5/2^+$	
				Γ (MeV)	$R = \Gamma(\Sigma_c^*\pi)/\Gamma(\Sigma_c\pi)$
Experimental values (Belle(07))				5.8 ± 1.1	$0.225 \pm 0.062 \pm 0.025$
	$(\ell_\lambda, \ell_\rho)$	$J_\Lambda^P(j)$	$(\ell_\lambda + \ell_\rho)$		
This work	(0, 1)	$5/2^-(2)$	1	42 – 55	1.6 – 1.7
	(2, 0)	$5/2^+(2)$	2	11 – 26	8.2 – 8.5
	(0, 2)	$5/2^+(2)$	2	28 – 52	19.0 – 19.1
	(1, 1)	$5/2^+(2)$	2	52 – 110	27.7 – 30.4
			1	0.63 – 1.7	(∞)
		$5/2^+(3)$	2	2.8 – 5.7	0.41 – 0.43

This is the only configuration to explain

Selection rule due to the brown muck (diquark)



Productions

- How much $\begin{cases} \text{charm} \\ \text{excited states} \end{cases}$ are produced
- Can we study structure?

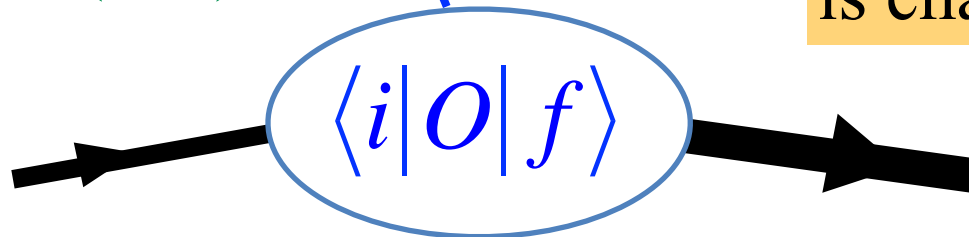
Productions

$$\pi + N \rightarrow D^* + \Lambda_c$$



A. B. Kaidalov and P. E. Volkovitsky,
B. Z. Phys. C 63, 517 (1994)

How much
is charm produced?



Quark model

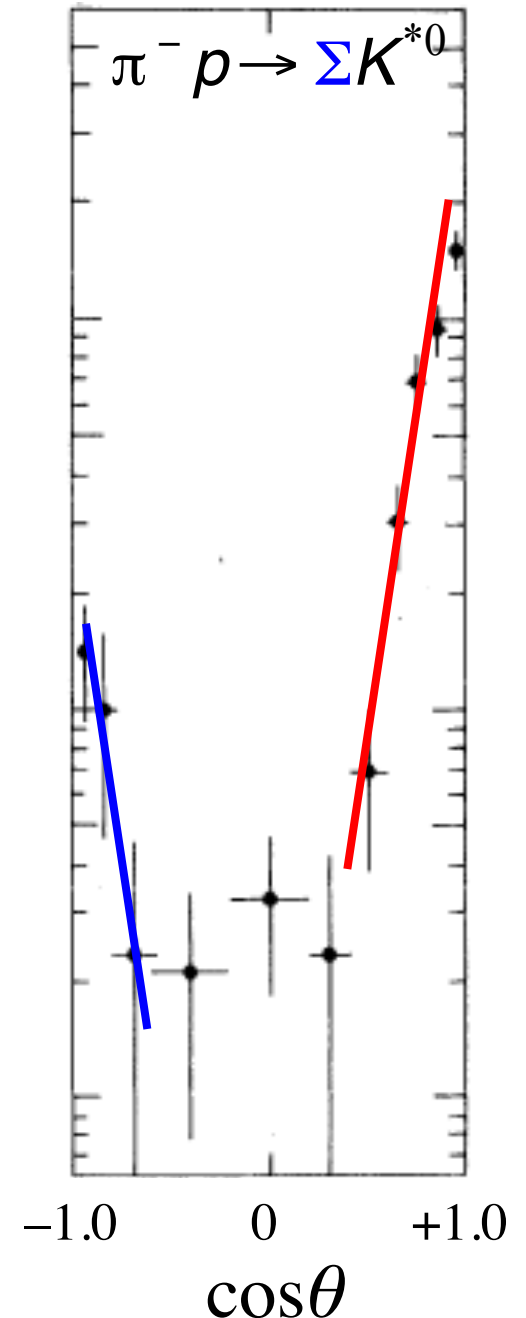
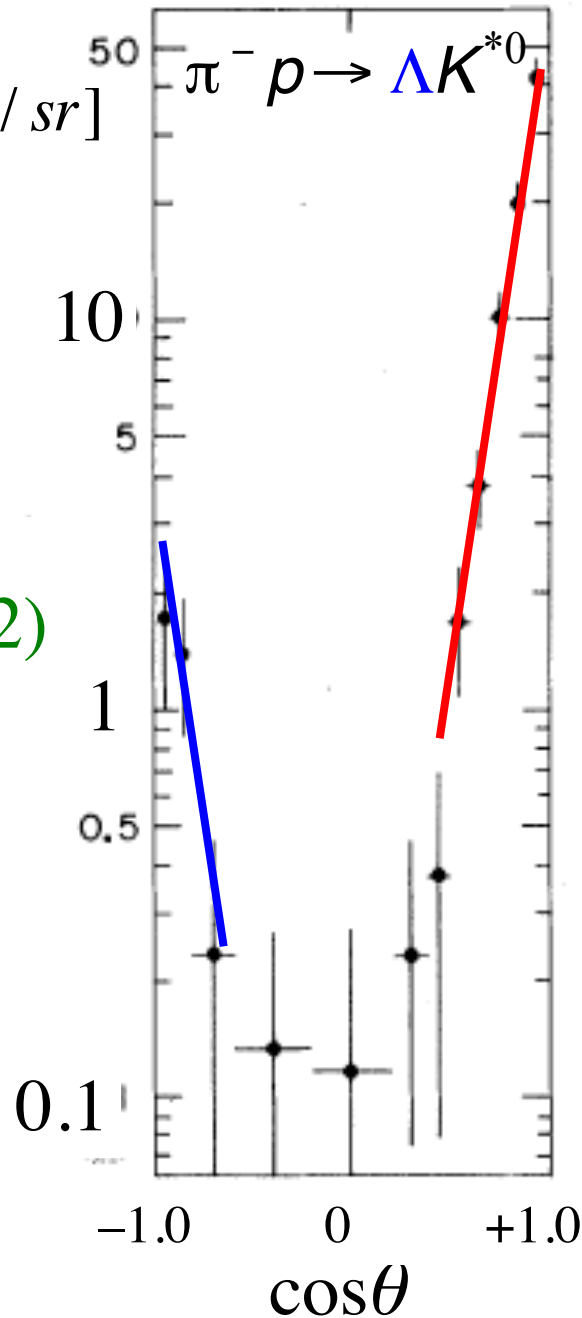
Kim, Kim, Noumi, Shirotori, Hosaka
PTEP 2014 (2014) 10, 103D01,
PRD92 (2015) 9, 094021

How are they related to
internal structure of Λ_c^* ?

$$\frac{d\sigma}{d\Omega} [\mu b / sr]$$

$$p_{\pi, \text{Lab}} = 4.5 \text{ GeV}$$

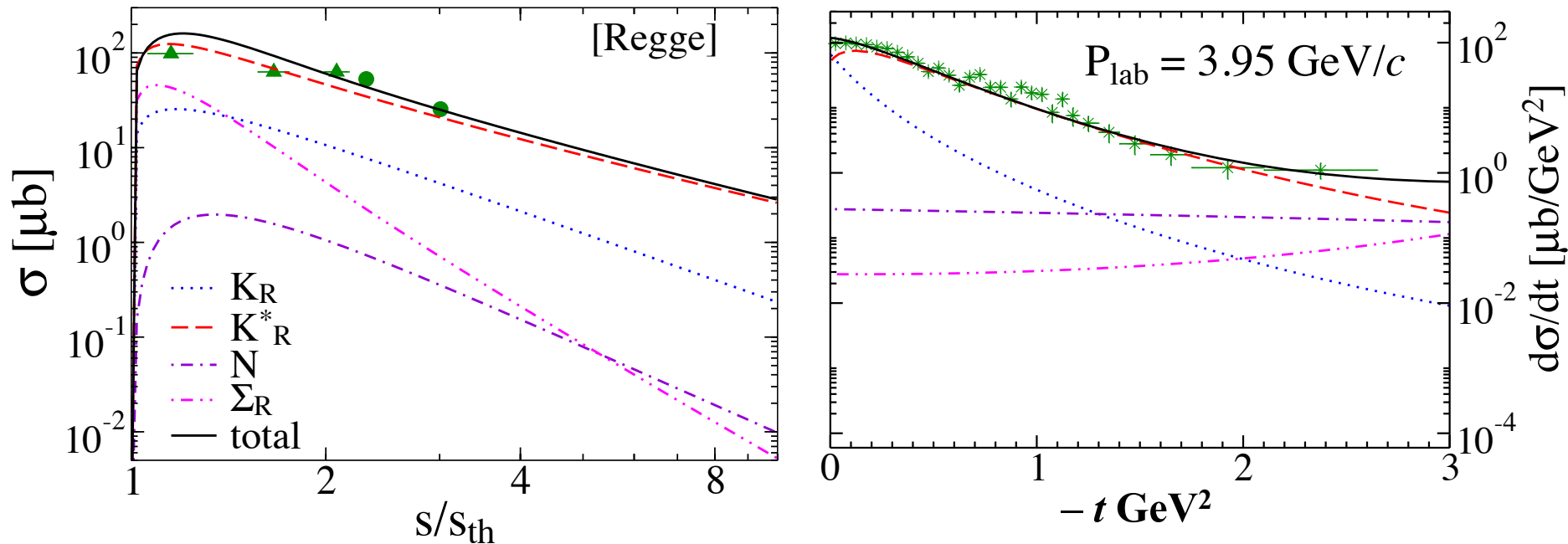
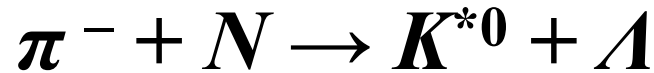
D.J. Krennel et al
PRD6, 1220 (1972)



How much is charm produced?

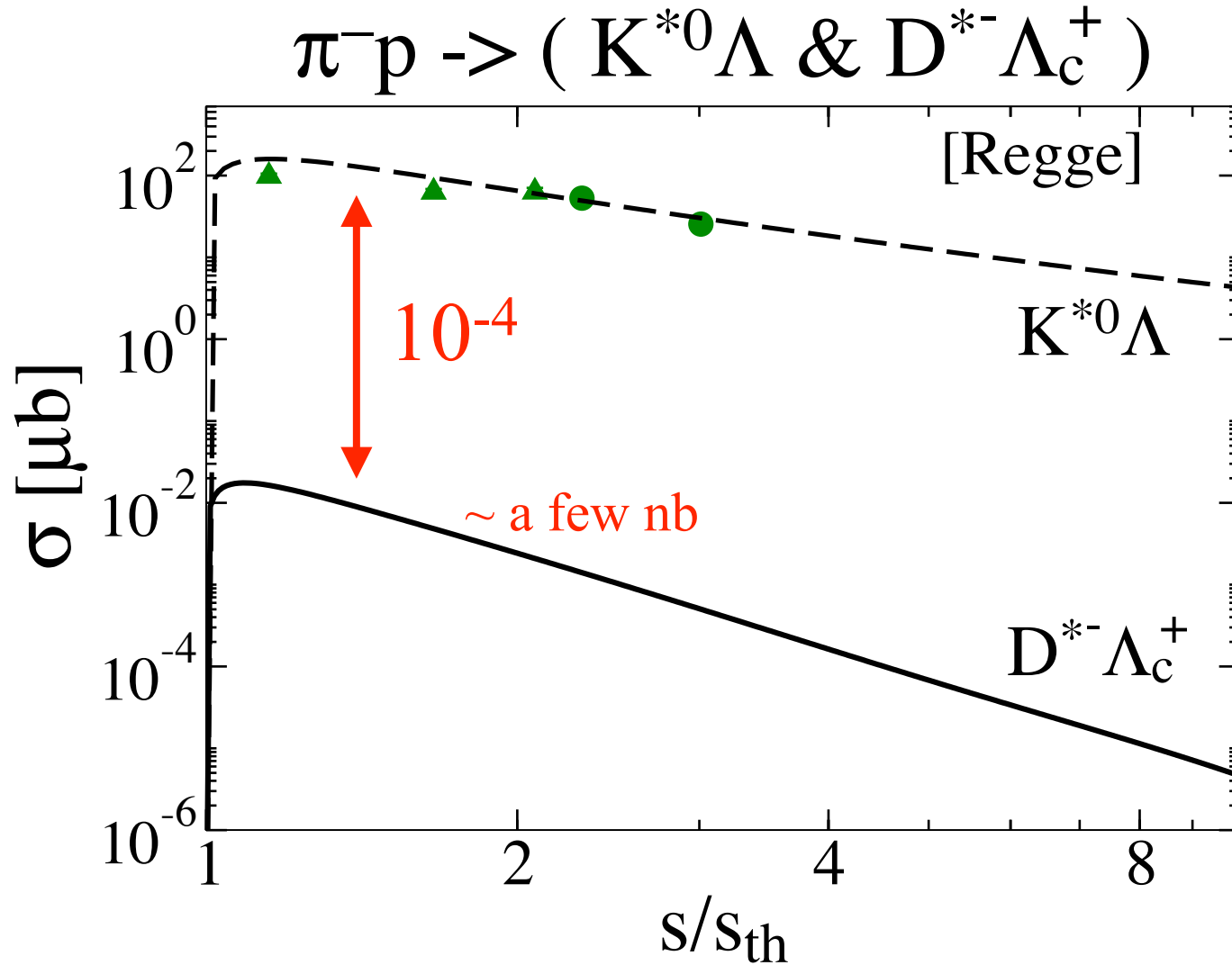
Regge model description

Kim Hosaka Kim Noumi, arXiv:1509.03567



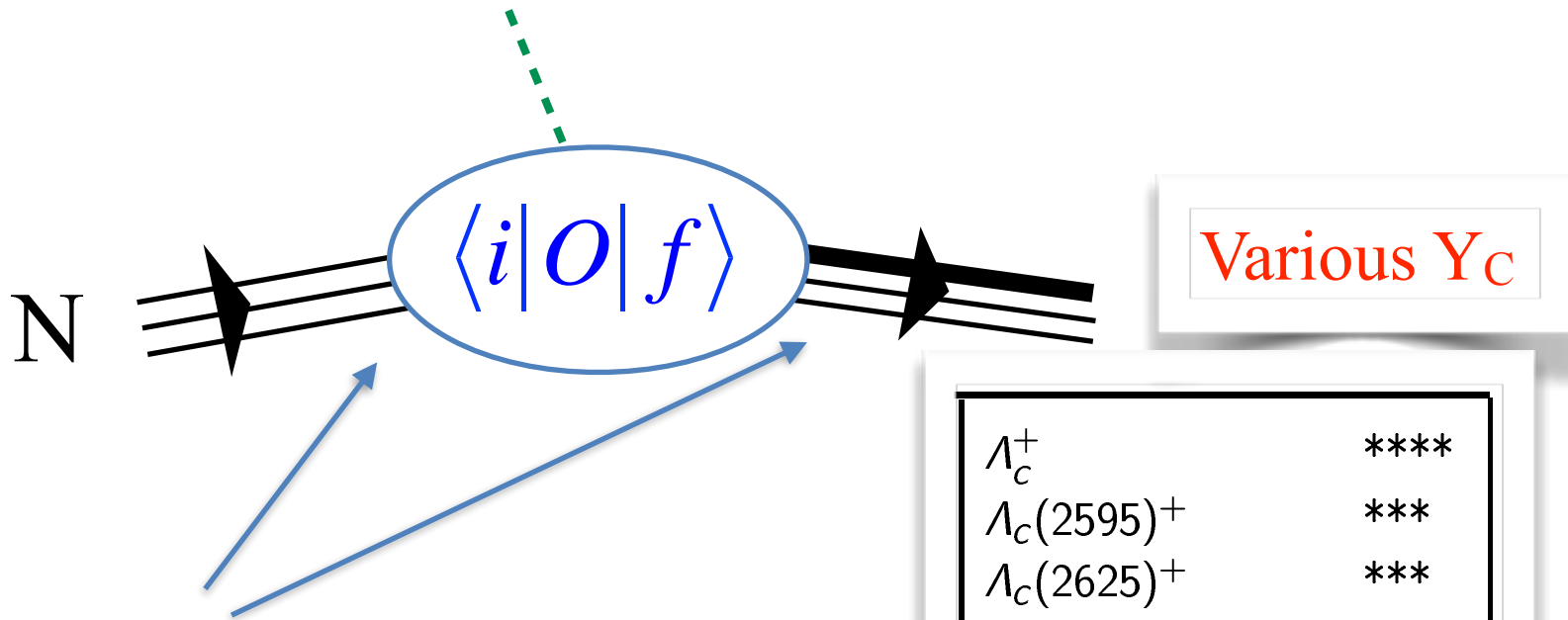
- Vector-Reggeon dominance with some pseudoscalar
- Energy dependence is also well produced

Prediction to the charm production



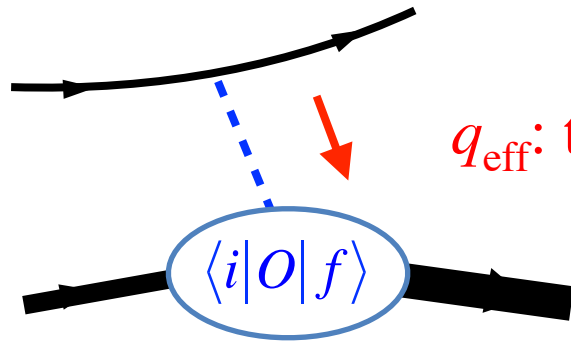
How are they related to internal structure?

PTEP 2014 (2014) 10, 103D01



Quark model wave functions
(Harmonic oscillator)

Λ_c^+	****
$\Lambda_c(2595)^+$	***
$\Lambda_c(2625)^+$	***
$\Lambda_c(2765)^+$	*
$\Lambda_c(2880)^+$	***
$\Lambda_c(2940)^+$	***
$\Sigma_c(2455)$	****
$\Sigma_c(2520)$	***
$\Sigma_c(2800)$	***



q_{eff} : the momentum transfer \sim Large

$$\text{GS} \quad \langle B_c(\text{S-wave}) | \vec{e}_\perp \cdot \vec{\sigma} e^{i\vec{q}_{\text{eff}} \cdot \vec{x}} | N(\text{S-wave}) \rangle_{\text{radial}} \sim 1 \times \exp\left(-\frac{q_{\text{eff}}^2}{4A^2}\right)$$

Excited states

$$\langle B_c(\text{P-wave}) | \vec{e}_\perp \cdot \vec{\sigma} e^{i\vec{q}_{\text{eff}} \cdot \vec{x}} | N(\text{S-wave}) \rangle_{\text{radial}} \sim \left(\frac{q_{\text{eff}}}{A}\right)^1 \times \exp\left(-\frac{q_{\text{eff}}^2}{4A^2}\right)$$

$$\langle B_c(\text{D-wave}) | \vec{e}_\perp \cdot \vec{\sigma} e^{i\vec{q}_{\text{eff}} \cdot \vec{x}} | N(\text{S-wave}) \rangle_{\text{radial}} \sim \left(\frac{q_{\text{eff}}}{A}\right)^2 \times \exp\left(-\frac{q_{\text{eff}}^2}{4A^2}\right)$$

Transitions to excited states are not suppressed!

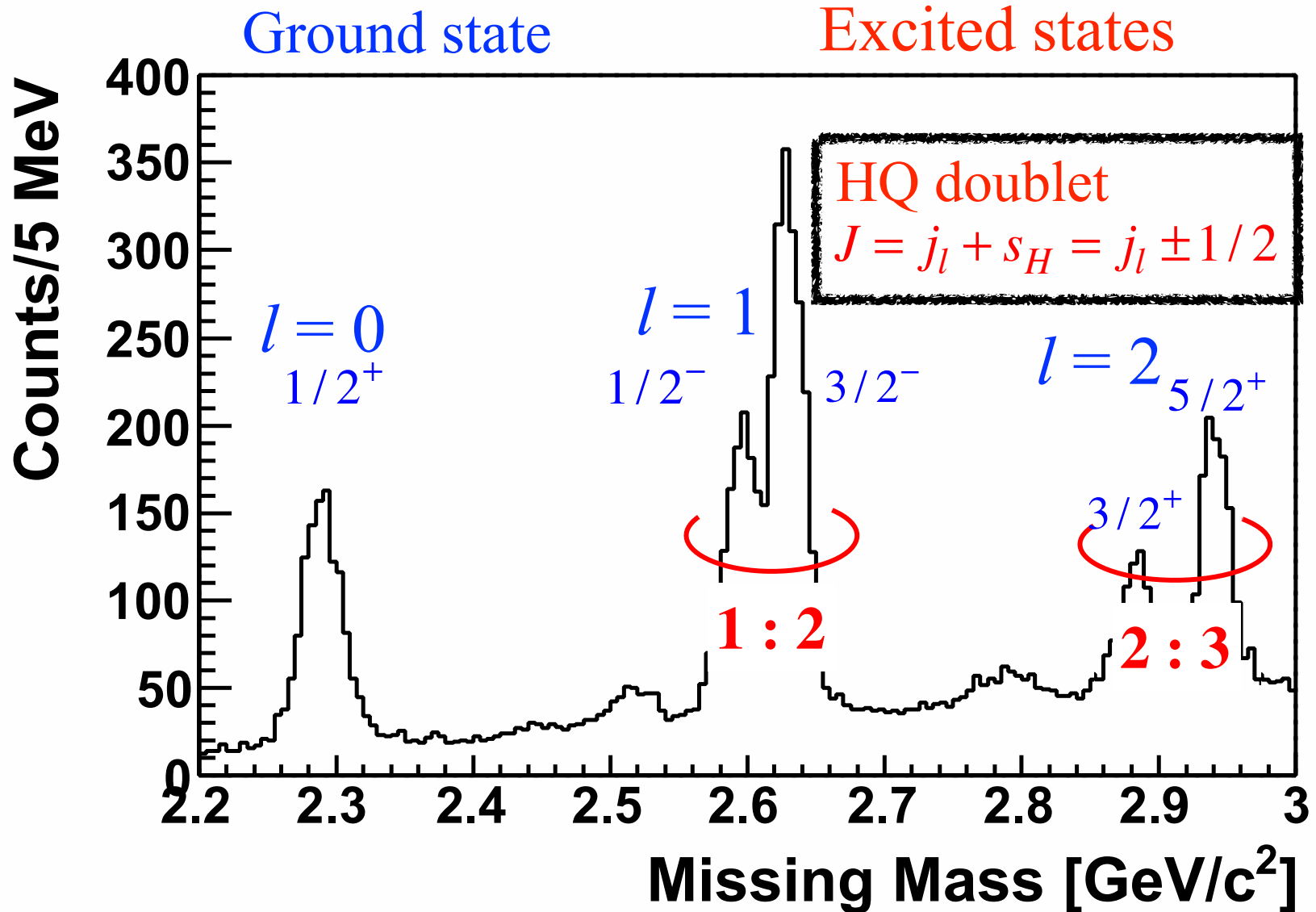
Charm $k_{\pi}^{CM} = 2.71$ [GeV], $k_{\pi}^{Lab} = 16$ [GeV]

$l = 0$	$\Lambda_c(\frac{1}{2}^+)$ 1.00	$\Sigma_c(\frac{1}{2}^+)$ 0.02	$\Sigma_c(\frac{3}{2}^+)$ 0.16						
<u>$l = 1$</u>	$\Lambda_c(\frac{1}{2}^-)$ 0.90	$\Lambda_c(\frac{3}{2}^-)$ 1.70	$\Sigma_c(\frac{1}{2}^-)$ 0.02	$\Sigma_c(\frac{3}{2}^-)$ 0.03	$\Sigma'_c(\frac{1}{2}^-)$ 0.04	$\Sigma'_c(\frac{3}{2}^-)$ 0.19	$\Sigma'_c(\frac{5}{2}^-)$ 0.18		
<u>$l = 2$</u>	$\Lambda_c(\frac{3}{2}^+)$ 0.50	$\Lambda_c(\frac{5}{2}^+ -)$ 0.88	$\Sigma_c(\frac{3}{2}^+)$ 0.02	$\Sigma_c(\frac{5}{2}^+)$ 0.02	$\Sigma'_c(\frac{1}{2}^+)$ 0.01	$\Sigma'_c(\frac{3}{2}^+)$ 0.03	$\Sigma'_c(\frac{5}{2}^+)$ 0.07	$\Sigma'_c(\frac{5}{2}^+)$ 0.07	

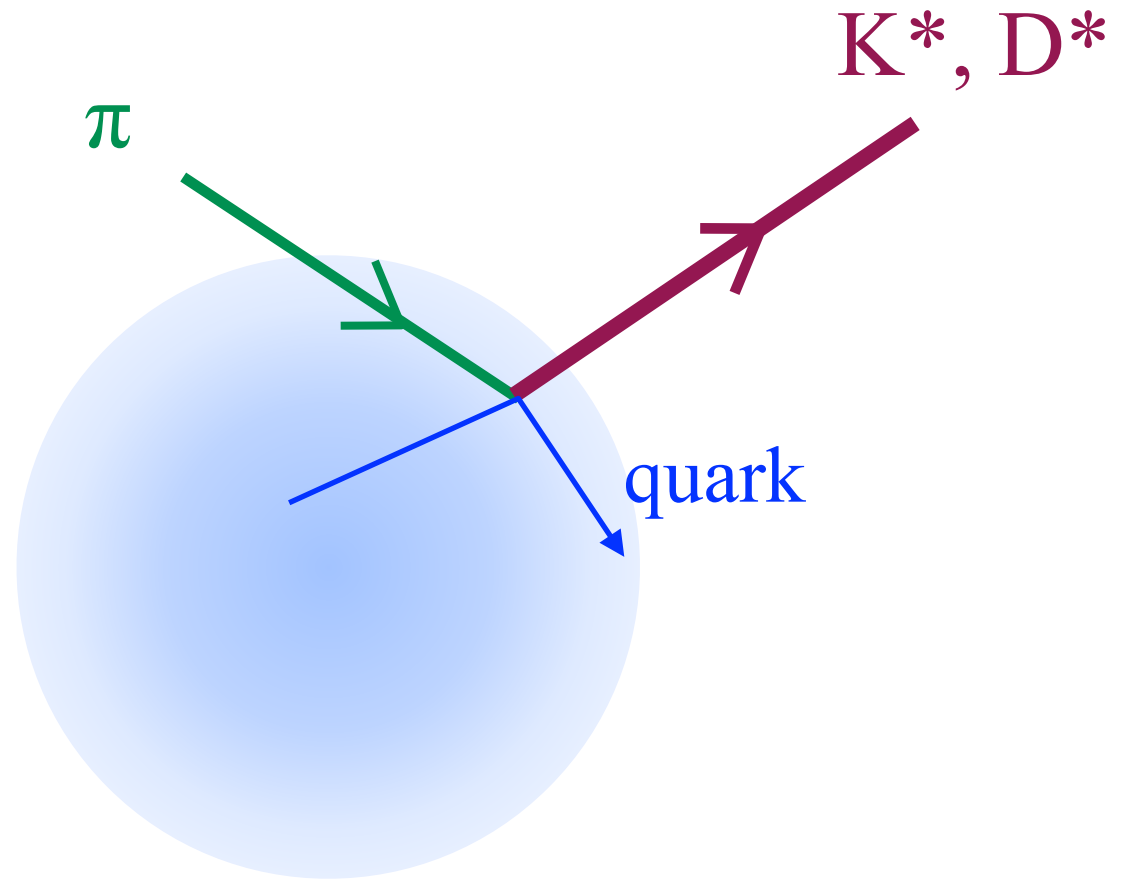
Strange $k_{\pi}^{CM} = 1.59$ [GeV], $k_{\pi}^{Lab} = 5.8$ [GeV]

$l = 0$	$\Lambda_-(\frac{1}{2}^+)$ 1.00	$\Sigma_-(\frac{1}{2}^+)$ 0.067	$\Sigma_-(\frac{3}{2}^+)$ 0.44						
$l = 1$	$\Lambda_-(\frac{1}{2}^-)$ 0.11	$\Lambda_-(\frac{3}{2}^-)$ 0.23	$\Sigma_-(\frac{1}{2}^-)$ 0.007	$\Sigma_-(\frac{3}{2}^-)$ 0.01	$\Sigma'_-(\frac{1}{2}^-)$ 0.01	$\Sigma'_-(\frac{3}{2}^-)$ 0.07	$\Sigma'_-(\frac{5}{2}^-)$ 0.067		
$l = 2$	$\Lambda_-(\frac{3}{2}^+)$ 0.13	$\Lambda_-(\frac{5}{2}^+ -)$ 0.20	$\Sigma_-(\frac{3}{2}^+)$ 0.007	$\Sigma_-(\frac{5}{2}^+)$ 0.01	$\Sigma'_-(\frac{1}{2}^+)$ 0.004	$\Sigma'_-(\frac{3}{2}^+)$ 0.02	$\Sigma'_-(\frac{5}{2}^+)$ 0.038	$\Sigma'_-(\frac{5}{2}^+)$ 0.04	

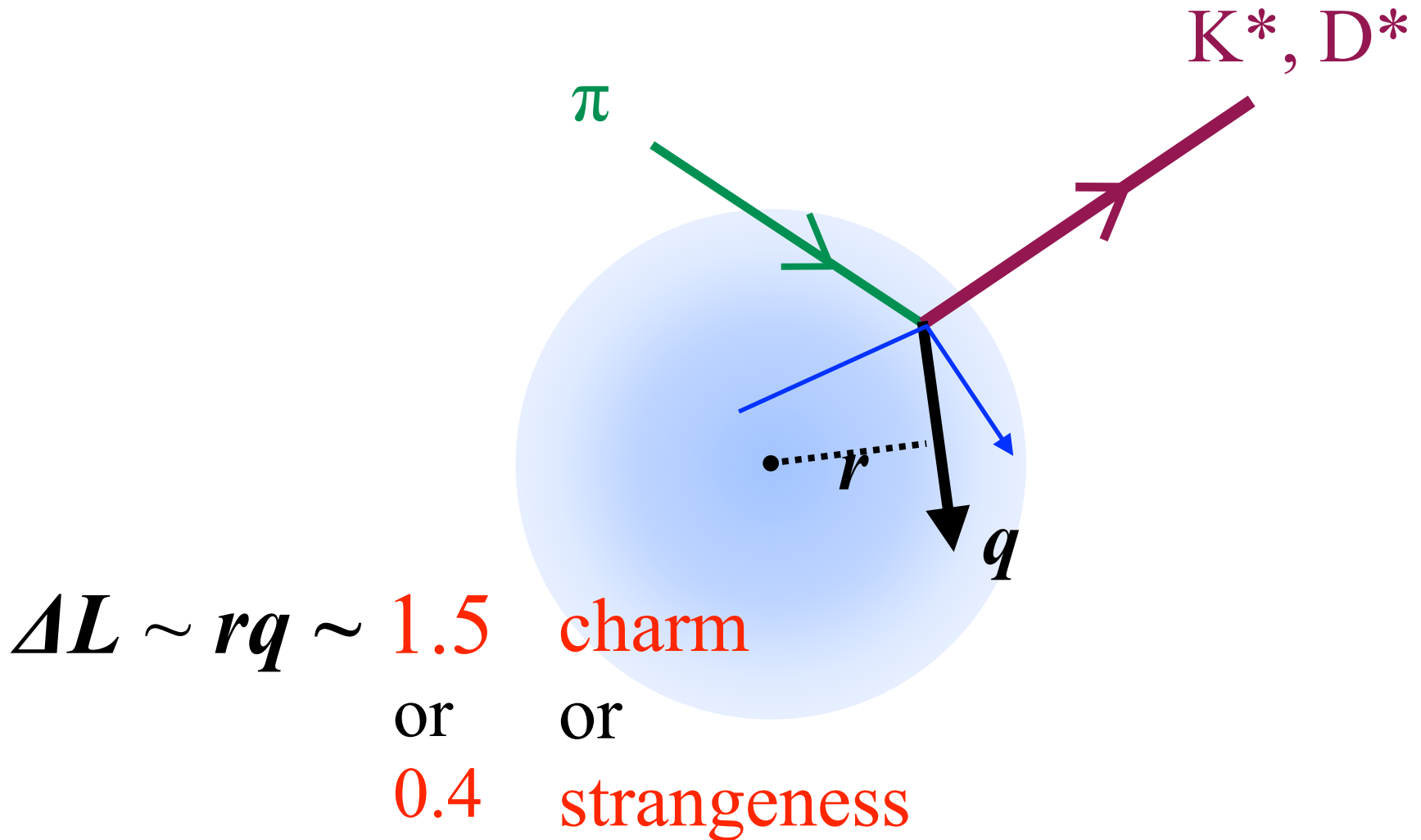
Charm production spectrum



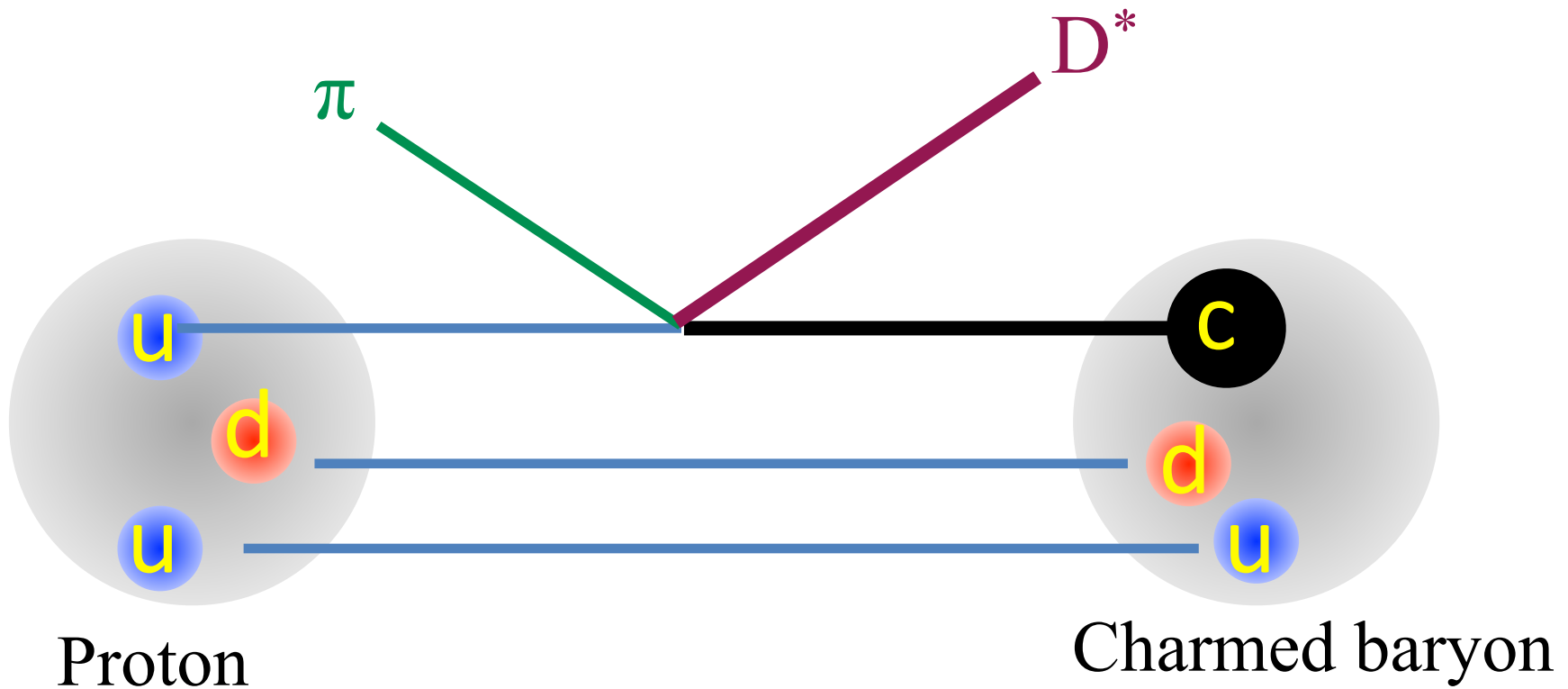
Angular momentum dependence



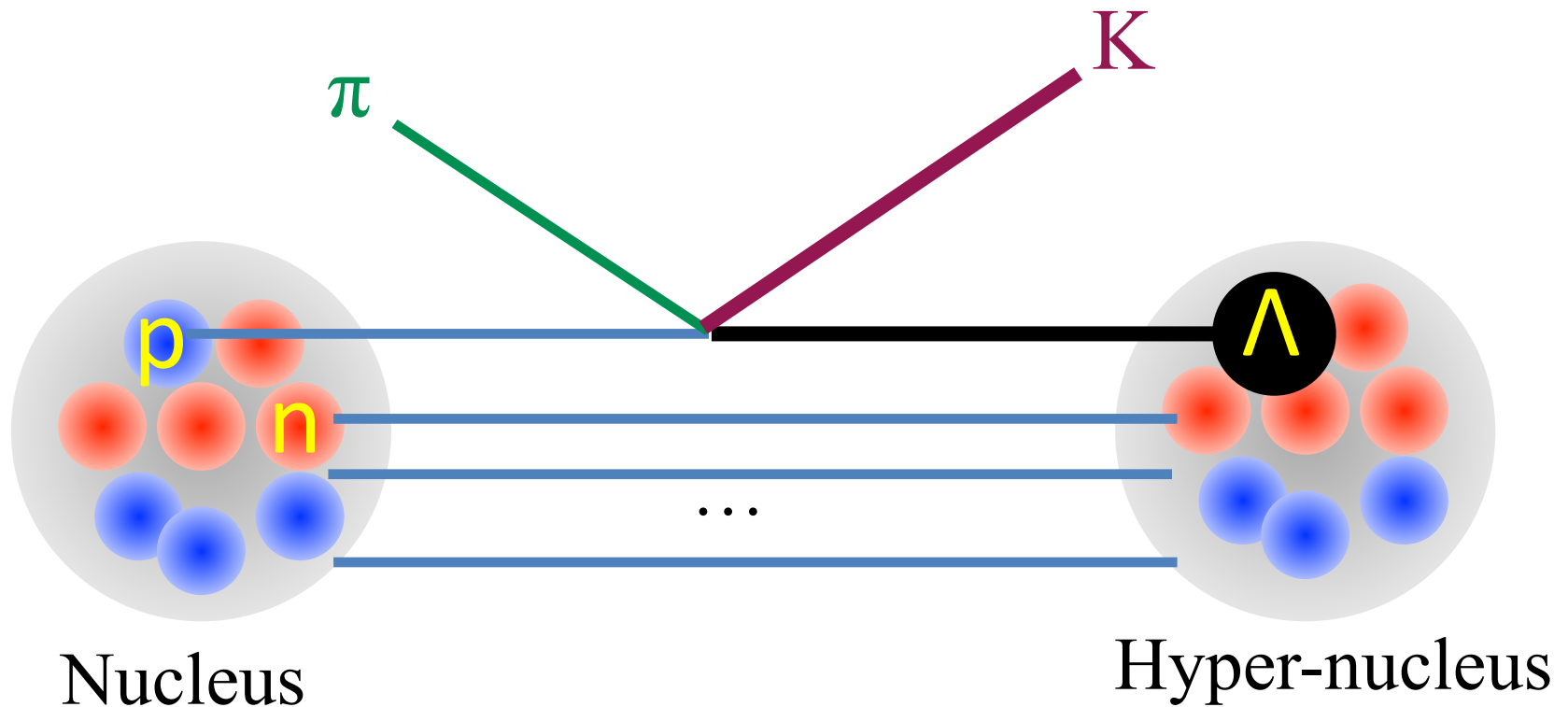
Angular momentum dependence



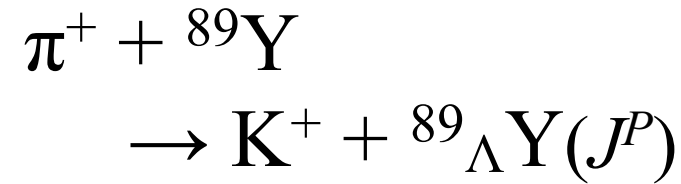
Analogy to the hyper nucleus production



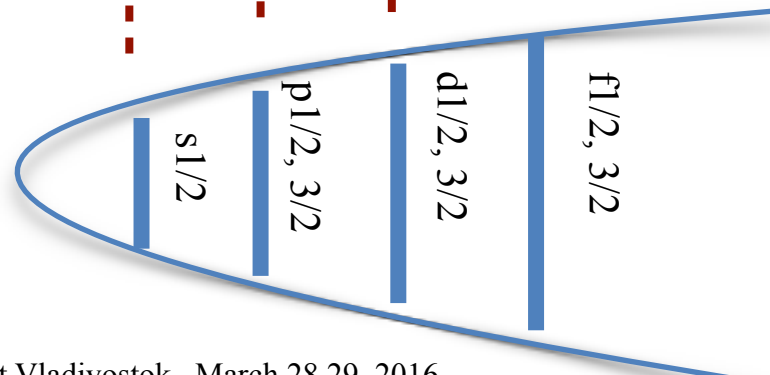
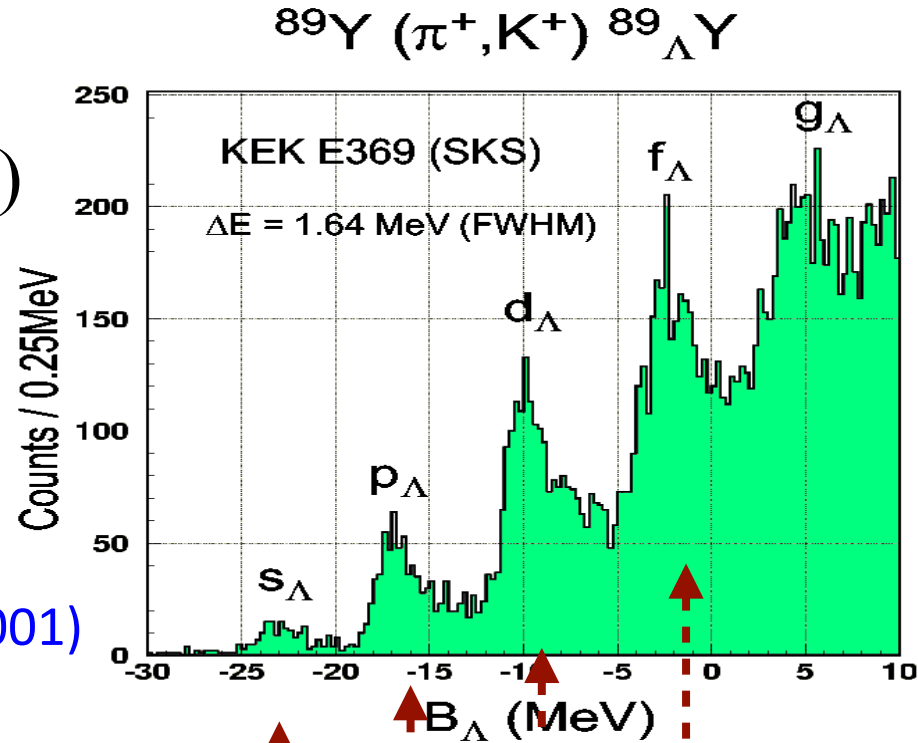
Analogy to the hyper nucleus production

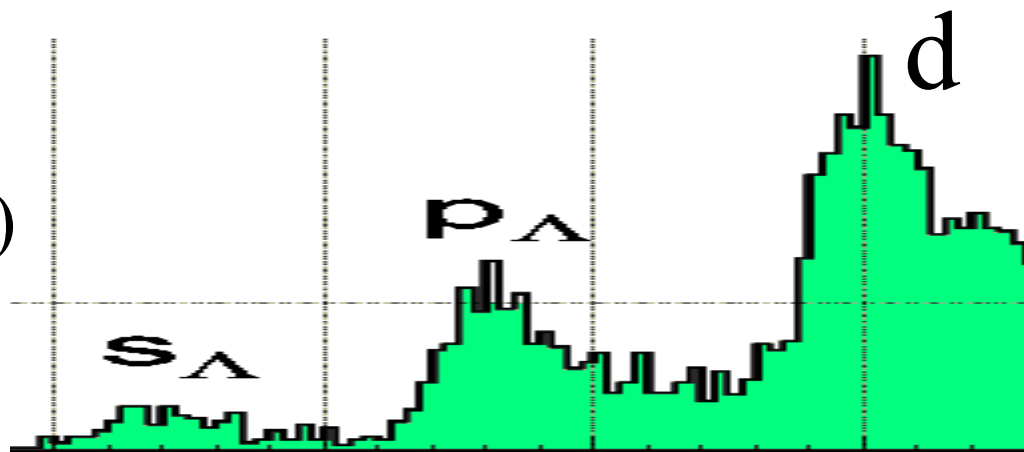
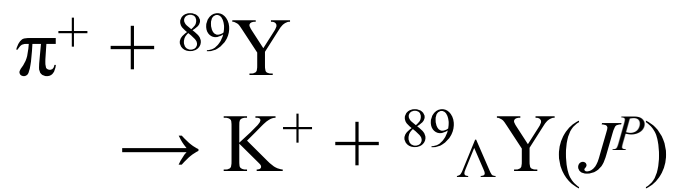


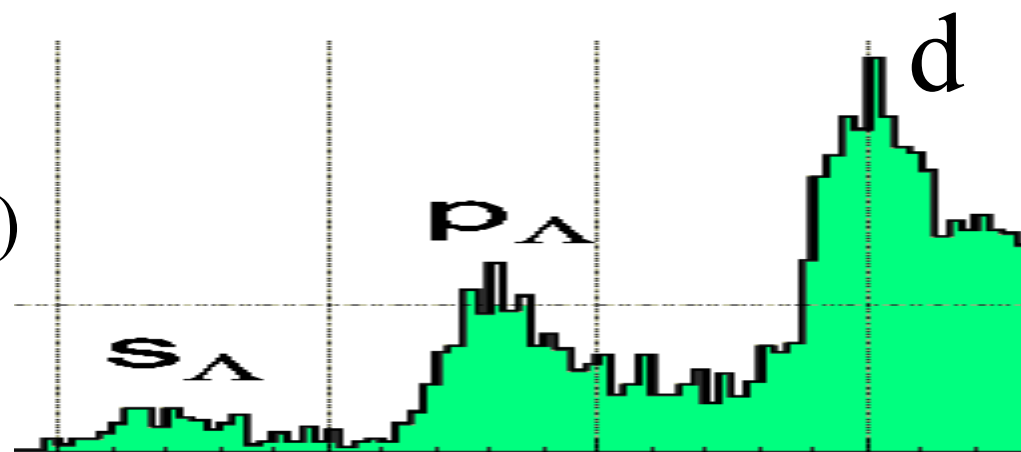
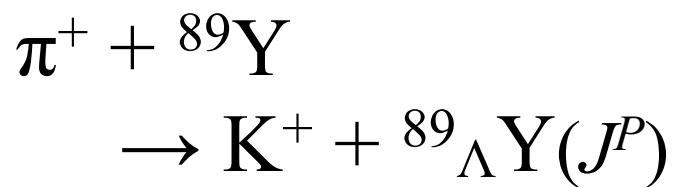
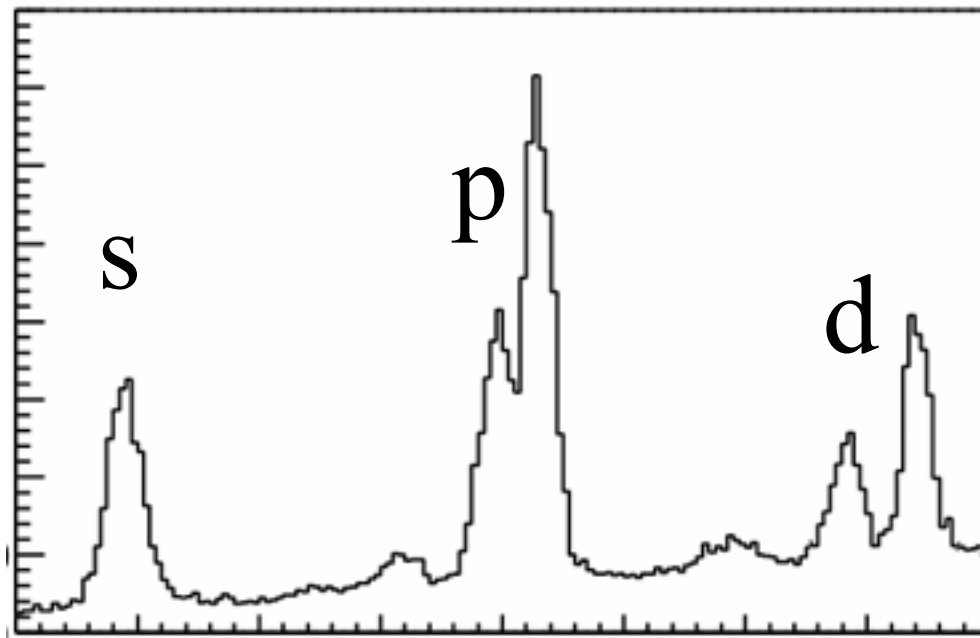
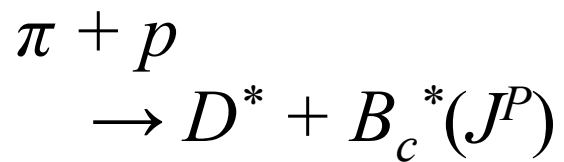
Analogy to the hyper nucleus production

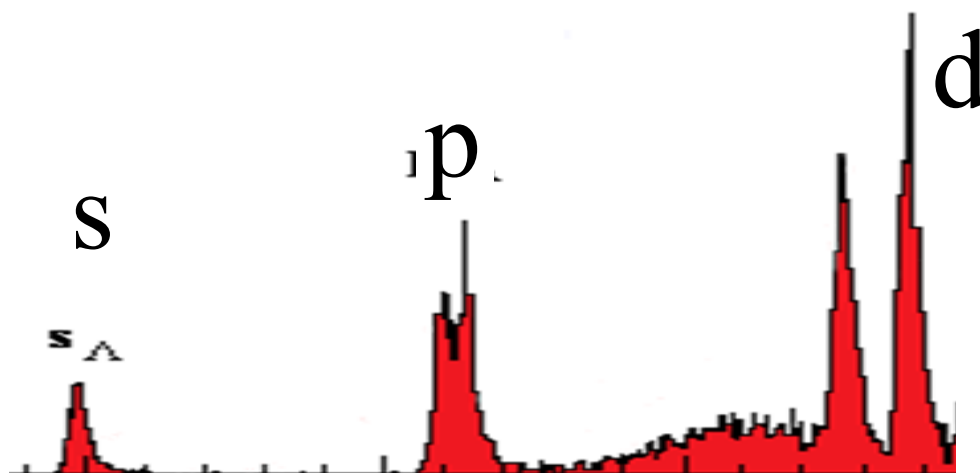
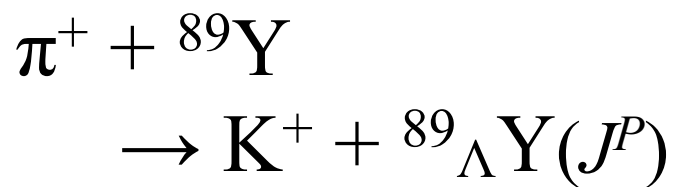
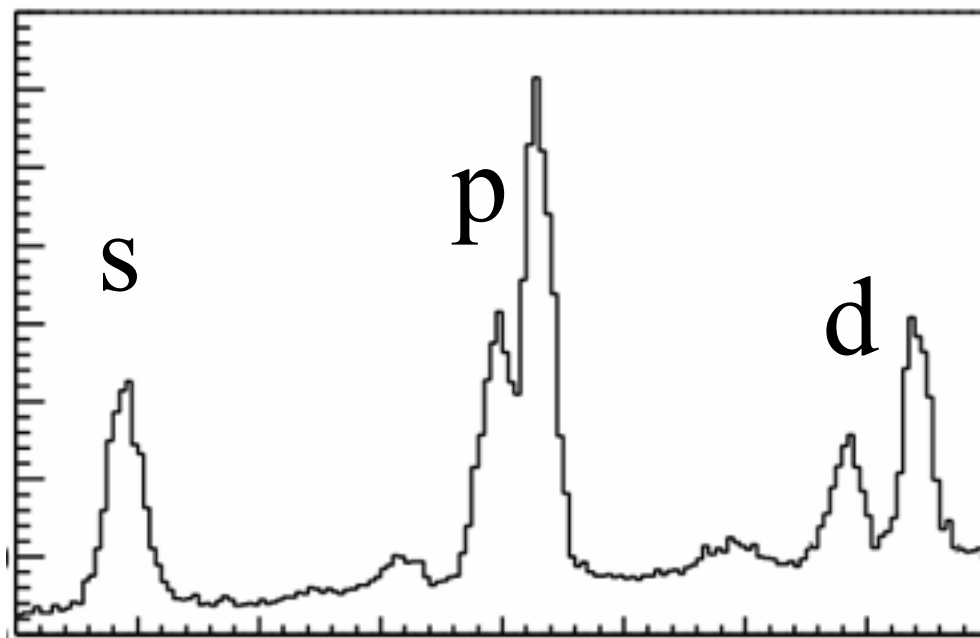
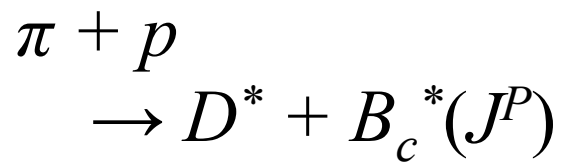


H. Hotch *et al.*,
Phys. Rev. C64, 044302(2001)









Summary

- Heavy quarks identify and disentangle different modes of baryons, ρ and λ modes \Rightarrow diquark dynamics?
- **Decays** are useful to further understand the structure and fundamental nature of hadron physics = QCD
- **Productions** are useful for structure study
A similar feature with hyper nuclei
- Charm baryons could be abundantly produced

Further studies at J-PARC